

**THERMAL STRESSES IN ANISOTROPIC PLATES**

**A Thesis**

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15

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## TABLE OF CONTENTS

	Page
ACKNOWLEDGEMENT . . . . .	ii
LIST OF FIGURES . . . . .	iv
NOTATION . . . . .	v
ABSTRACT . . . . .	vi
<b>Chapter</b>	
I. INTRODUCTION . . . . .	1
II. THERMOELASTIC EQUATIONS FOR ANISOTROPIC PLATES . . . . .	4
III. SOLUTION APPROACH . . . . .	15
Finite Difference Equation for F . . . . .	17
Finite Difference Equation for W . . . . .	20
Boundary Conditions--All Edges Free . . . . .	23
Boundary Conditions in Terms of F . . . . .	23
Boundary Conditions in Terms of W . . . . .	24
IV. THE COMPUTER PROGRAM . . . . .	29
RESULTS AND CONCLUSIONS . . . . .	36
BIBLIOGRAPHY . . . . .	37
APPENDIX A . . . . .	38
APPENDIX B . . . . .	62
VITA . . . . .	74

LIST OF FIGURES

Figure	Page
1. Coordinate System and the Plate Dimensions . . . . .	4
2. A Grid Superimposed on a Plate . . . . .	15
3. Moments and Shear Forces Acting on a Volume Element . . . . .	25
4. Main Program Flow Chart . . . . .	31
5. Subroutine FUNCF Flow Chart . . . . .	32
6. Flow Charts for Subroutines TEMDIS (6a) and SIGMA0 (6b) . . . . .	33
7. Subroutine DEFLXN Flow Chart . . . . .	34
8. Subroutine SIGMA1 Flow Chart . . . . .	35

## NOTATION

$T, T_0, T_1$	Temperature distributions in the plate
$a, b, 2h$	Length, width, and thickness of the plate, respectively
$u, v$	Displacements of a point $(X, Y, Z)$ of the plate in X and Y directions, respectively
$w$	Deflection of middle surface of the plate
$\tau, \tau^0, \tau^1$	Stress tensors corresponding to $T, T_0, T_1$ distribution
$\sigma_x, \sigma_y, \sigma_z$	Average normal stress in direction of the axis indicated by subscript
$N_x, N_y$	Normal forces per unit of length in X and Y directions, respectively
$N_{xy}, Q_x, Q_y$	Shear forces per unit of length in the planes XY, XZ, and YZ, respectively
$a_{ij}, c_{ij}$	Elastic constants (compliances and stiffnesses, respectively)
$b_{ij}$	Modified stiffnesses
$\alpha_i$	Linear coefficients of thermal expansion
$\alpha_i^1$	Modified linear coefficients of thermal expansion
$\epsilon_x, \epsilon_y, \gamma_{xy}$	Unit elongations in the X and Y directions, and shearing strain component in the XY plane, respectively
$F$	Stress function
$\Delta h, \Delta k$	Finite increments in the X and Y directions, respectively, for solving finite difference equations
$N_x, N_y$	Number of grid points in the X and Y directions, respectively
$M_x, M_y, M_{xy}$	Moments per unit of length of the stresses $\sigma_x, \sigma_y, \tau_{xy}$ about the middle surface of the plate, respectively.

## ABSTRACT

Many materials used in the nuclear industry have anisotropic properties. This thesis investigates stresses within a simple body due to temperature variations. A computer program is presented which can be used to calculate thermal stresses and deflections in an anisotropic plate with all four edges free and having one plane of elastic symmetry. General fourth order polynomials are assumed to predict the correct temperature distributions. Stresses are first calculated for a two-dimensional (X and Y) temperature distribution with no deflection (in Z direction) allowed. Then a three-dimensional case with a linear temperature gradient in Z direction is considered. Superimposition of these solutions leads to a general treatment of the normal and shear stresses in the plate. The correct elastic constants for the material being considered are required as input to the program.

## CHAPTER I

### INTRODUCTION

Anisotropic materials play an important part in modern technology. These find applications in a number of areas such as the aerospace, missile, aircraft and nuclear industries. It is possible to produce an artificial anisotropy in isotropic materials. For example, corrugated plates and membranes can be made from elastically isotropic materials. Certain structures which are strengthened by ribbing also become anisotropic.

An elastic medium is isotropic if in all directions its elastic properties at all points are the same. It will be anisotropic if its elastic properties at any point in the medium are different in different directions. An anisotropic body may have one, two or three orthogonal planes of elastic symmetry passing through any point in the body. The plane of elastic symmetry is the plane of the body about which its elastic properties are symmetric. The body which has three orthogonal planes of elastic symmetry at each point is called orthogonally anisotropic, or simply, orthotropic. If there are sets of equivalent directions in a body which can be superimposed by a rotation of  $2\pi/N$  about a certain axis, say Z-axis, then this is an axis of elastic symmetry of order N. A plane normal to an axis of elastic symmetry of an infinitely large order ( $N \rightarrow \infty$ ) is a plane of isotropy. A body having this property is called a transversely anisotropic body.

This is also described as a transversely isotropic body.

Under certain load conditions, such as thermal or external forces, an elastic body undergoes deformations or shape changes. Stresses, that is, loads per unit area, are related to these deformations by the well known Hooke's law for elastic materials. In the case where there is no motion of the body upon application of the forces, the body as a whole must be in equilibrium. If a small volume element in the body is considered, certain relations known as equilibrium equations arise. Also, any body has certain constraints acting on its boundaries. Using Hooke's law, the equilibrium equations and the boundary conditions simultaneously could lead to a solution for the distribution of the stresses in the body.

Although the linear theory of elasticity for an anisotropic medium has been known for quite a long time, a realistic solution has not been presented except in the case of orthotropic materials [1]. Solution has been presented for an anisotropic plate with one plane of elastic symmetry [2] but it involves a number of mappings into different planes and does not reach a final expression. Equations for the deflection of an isotropic thin elastic plate subjected to a temperature distribution of the form

$$T(x,y,z) = T_0(x,y) + zT_1(x,y) \quad (1)$$

have been derived by Nadai [3]. He did not, however, work on the solutions of these equations. Principally, the isothermal theory of anisotropic plates has been developed by Boussinesq [4], Voigt [5] and Lechnitzky [6]. It appears from a literature survey that until 1941



only Voigt [5] had given any consideration to the thermal effects on the anisotropic plate. He, too, considered a simple case with temperature distribution of the kind

$$T = T(x,y)$$

and, thus, did not consider any bending effect. In 1941, Sokolonikoff [7] presented a short summary of anisotropic plate theory. Pell [2] considered the general case of a plate acted upon by forces on its edges, free of body forces and having a temperature distribution as given by equation (1). After using a number of mapping functions, he arrived at a rather complicated solution.

It is the aim of this research to develop a computer program to compute the thermal deflections and stresses in an anisotropic plate that has at least one plane of elastic symmetry. The temperature distribution considered is given by equation (1). Introduction of a stress function  $F$  reduces the compatibility equation to a fourth order partial differential equation [7]. The statical equilibrium equation reduces to a fourth order partial differential equation in terms of  $W$ , the deflection. These two equations are solved on a digital computer using the finite difference technique. Once the values of  $F$  have been found, the stress distribution due to  $T_0$  can be obtained by direct substitution of  $F$  into the corresponding stress equations. In a similar way, the values of  $W$  lead to stress distribution due to  $zT_1$ . Superposition of the latter onto the earlier ones gives the total stress distribution in the plate.

## CHAPTER II

### THERMOELASTIC EQUATIONS FOR ANISOTROPIC PLATES

Consider a thin elastic plate composed of a medium having at least one plane of elastic symmetry passing through each point parallel to the middle plane of the plate, which is chosen to lie in the  $x$ - $y$  plane. Let the plate be subjected to a temperature distribution given by equation (1). The Cartesian coordinate system and the plate dimensions are shown in Figure 1.

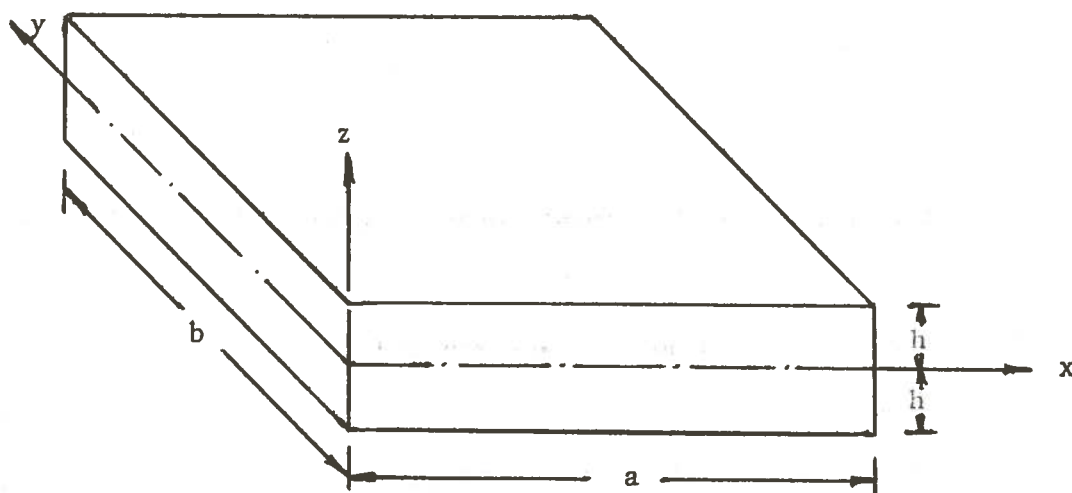


Figure 1

#### Coordinate System and the Plate Dimensions

The Kirchhoff assumptions [7] of the thin plate theory lead to the fundamental relations

$$u = -z \frac{\partial w}{\partial x} , \quad (2)$$

$$v = -z \frac{\partial w}{\partial y} ,$$

and

$$|\sigma_z| \ll \{ |\sigma_x|, |\sigma_y|, |\tau_{xy}| \} , \quad (3)$$

where  $\sigma_x$ ,  $\sigma_y$ ,  $\sigma_z$  are the normal stresses in x, y and z directions, respectively, and  $\tau_{xy}$  is the shear stress in the x,y plane. In these equations w is the deflection of the middle surface of the plate, and u and v are the displacements of a point (x,y,z) of the plate in the x and y directions respectively. These relations are valid throughout the thickness, 2h, of the plate. Pell [2] further assumed that the stress tensor  $\tau$  at any point in the plate is the linear sum of two stress tensors:

$$\tau = \tau^0 + \tau^1 , \quad (4)$$

where  $\tau^0$  is a plane stress tensor generated by  $T_0(x,y)$ , and  $\tau^1$  is the bending stress tensor in the plate due to  $zT_1$ .

The generalized Hooke's law [8] connecting the components of stress at any point of an elastic solid body with the components of strain at that point, taking into account the thermal effects, is expressed by

$$\frac{\partial u}{\partial x} = a_{11}\sigma_x + a_{12}\sigma_y + a_{13}\sigma_z + a_{16}\tau_{xy} + \alpha_1 T,$$

$$\frac{\partial v}{\partial y} = a_{21}\sigma_x + a_{22}\sigma_y + a_{23}\sigma_z + a_{26}\tau_{xy} + \alpha_2 T,$$

$$\frac{\partial w}{\partial z} = a_{31}\sigma_x + a_{32}\sigma_y + a_{33}\sigma_z + a_{36}\tau_{xy} + \alpha_3 T,$$

(5)

$$\frac{\partial v}{\partial z} + \frac{\partial w}{\partial y} = a_{44}\tau_{yz} + a_{45}\tau_{xz} + 2\alpha_4 T,$$

$$\frac{\partial w}{\partial x} + \frac{\partial u}{\partial z} = a_{54}\tau_{yz} + a_{55}\tau_{xz} + 2\alpha_5 T,$$

and

$$\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} = a_{61}\sigma_x + a_{62}\sigma_y + a_{63}\sigma_z + a_{66}\tau_{xy} + 2\alpha_6 T,$$

where the  $a_{ij}$  are elastic constants generally called compliances, and  $\alpha_i$  are the linear coefficients of thermal expansion of the medium.

Love [8] established the relationship

$$a_{ij} = a_{ji}$$

which helps reduce the number of compliances in equations (5) from 20 to 13.

For the plane stress condition, the system of equations (5) reduces to

$$\frac{\partial u^0}{\partial x} = a_{11}\sigma_x^0 + a_{12}\sigma_y^0 + a_{16}\tau_{xy}^0 + \alpha_1 T_0,$$

$$\frac{\partial v^0}{\partial y} = a_{12}\sigma_x^0 + a_{22}\sigma_y^0 + a_{26}\tau_{xy}^0 + \alpha_2 T_0,$$

(6)

and

$$\frac{\partial u^0}{\partial y} + \frac{\partial v^0}{\partial x} = a_{16}\sigma_x^0 + a_{26}\sigma_y^0 + a_{66}\tau_{xy}^0 + 2\alpha_6 T_0.$$

In solving two-dimensional problems by means of a stress formulation, it is convenient to introduce a stress function defined in such a way that the equations of equilibrium are identically satisfied.

*(other  
a<sub>ij</sub>'s drop  
out.)  
why they  
do*

In such a condition, the function will then be required to satisfy the compatibility equation [9]

$$\frac{\partial^2 \epsilon_x}{\partial y^2} + \frac{\partial^2 \epsilon_y}{\partial x^2} = 2 \frac{\partial^2 \gamma_{xy}}{\partial x \partial y} , \quad (7)$$

where  $\epsilon_x$  and  $\epsilon_y$  are unit elongations in the x and y directions and  $\gamma_{xy}$  is shearing strain component in the x,y plane.

In this case the stress function F is introduced in the following manner:

$$\sigma_x^0 = \frac{\partial^2 F}{\partial y^2} ; \sigma_y^0 = \frac{\partial^2 F}{\partial x^2} ; \tau_{xy}^0 = - \frac{\partial^2 F}{\partial x \partial y} . \quad (8)$$

It is further convenient in plate theory to deal with forces and moments per unit of length rather than with the stresses themselves. The following quantities are therefore introduced:

$$N_x = 2h\sigma_x^0 , N_y = 2h\sigma_y^0 , N_{xy} = 2h\tau_{xy}^0 , \quad (9)$$

where  $\sigma_x^0$  ,  $\sigma_y^0$  , and  $\tau_{xy}^0$  are the average normal stresses in the X and Y directions and the average shearing stress over the plate thickness, respectively.

From equations (6) and (8) one would obtain

$$\begin{aligned} \epsilon_x &= a_{11} \frac{\partial^2 F}{\partial y^2} + a_{12} \frac{\partial^2 F}{\partial x^2} - a_{16} \frac{\partial^2 F}{\partial x \partial y} + \alpha_1 T_0 , \\ \epsilon_y &= a_{12} \frac{\partial^2 F}{\partial y^2} + a_{22} \frac{\partial^2 F}{\partial x^2} - a_{26} \frac{\partial^2 F}{\partial x \partial y} + \alpha_2 T_0 , \end{aligned} \quad (10)$$

and

$$\gamma_{xy} = \frac{1}{2} \left\{ a_{16} \frac{\partial^2 F}{\partial y^2} + a_{26} \frac{\partial^2 F}{\partial x^2} - a_{66} \frac{\partial^2 F}{\partial x \partial y} + 2\alpha_6 T_0 \right\} .$$

Substituting these expressions into equation (7) results in

$$\begin{aligned}
 & a_{11} \frac{\partial^4 F}{\partial y^4} + a_{12} \frac{\partial^4 F}{\partial x^2 \partial y^2} - a_{16} \frac{\partial^4 F}{\partial x \partial y^3} + \alpha_1 \frac{\partial^2 T_0}{\partial y^2} + a_{12} \frac{\partial^4 F}{\partial x^2 \partial y^2} \\
 & + a_{22} \frac{\partial^4 F}{\partial x^4} - a_{26} \frac{\partial^4 F}{\partial x^3 \partial y} + \alpha_2 \frac{\partial^2 T_0}{\partial x^2} \\
 & = a_{16} \frac{\partial^4 F}{\partial x \partial y^3} + a_{26} \frac{\partial^4 F}{\partial x^3 \partial y} - a_{66} \frac{\partial^4 F}{\partial x^2 \partial y^2} + 2\alpha_6 \frac{\partial^2 T_0}{\partial x \partial y} .
 \end{aligned}$$

Rearranging this equation leads to the relation

$$\begin{aligned}
 & a_{22} \frac{\partial^4 F}{\partial x^4} - 2a_{26} \frac{\partial^4 F}{\partial x^3 \partial y} + (2a_{12} + a_{66}) \frac{\partial^4 F}{\partial x^2 \partial y^2} - 2a_{16} \frac{\partial^4 F}{\partial x \partial y^3} + a_{11} \frac{\partial^4 F}{\partial y^4} \\
 & = - \left\{ \alpha_2 \frac{\partial^2 T_0}{\partial x^2} - 2\alpha_6 \frac{\partial^2 T_0}{\partial x \partial y} + \alpha_1 \frac{\partial^2 T_0}{\partial y^2} \right\} . \quad (11)
 \end{aligned}$$

The generalized Hooke's law [8], taking into account the thermal effects due to  $zT_1$ , is expressed by

$$\begin{aligned}
 \sigma_x^1 &= c_{11} \left( \frac{\partial u^1}{\partial x} - \alpha_1 z T_1 \right) + c_{12} \left( \frac{\partial v^1}{\partial y} - \alpha_2 z T_1 \right) \\
 &+ c_{13} \left( \frac{\partial w^1}{\partial z} - \alpha_3 z T_1 \right) + c_{16} \left( \frac{1}{2} \left[ \frac{\partial u^1}{\partial y} + \frac{\partial v^1}{\partial x} \right] - \alpha_6 z T_1 \right) , \\
 \sigma_y^1 &= c_{12} \left( \frac{\partial u^1}{\partial x} - \alpha_1 z T_1 \right) + c_{22} \left( \frac{\partial v^1}{\partial y} - \alpha_2 z T_1 \right) \\
 &+ c_{23} \left( \frac{\partial w^1}{\partial z} - \alpha_3 z T_1 \right) + c_{26} \left( \frac{1}{2} \left[ \frac{\partial u^1}{\partial y} + \frac{\partial v^1}{\partial x} \right] - \alpha_6 z T_1 \right) ,
 \end{aligned} \quad (12)$$

$$\begin{aligned} \sigma_z^1 &= c_{13} \left( \frac{\partial u^1}{\partial x} - \alpha_1 z T_1 \right) + c_{23} \left( \frac{\partial v^1}{\partial y} - \alpha_2 z T_1 \right) \\ &+ c_{33} \left( \frac{\partial w^1}{\partial z} - \alpha_3 z T_1 \right) + c_{36} \left( \frac{1}{2} \left[ \frac{\partial u^1}{\partial y} + \frac{\partial v^1}{\partial x} \right] - \alpha_6 z T_1 \right), \end{aligned} \quad (12 \text{ cont'd})$$

and

$$\begin{aligned} \tau_{xy}^1 &= c_{16} \left( \frac{\partial u^1}{\partial x} - \alpha_1 z T_1 \right) + c_{26} \left( \frac{\partial v^1}{\partial y} - \alpha_2 z T_1 \right) \\ &+ c_{36} \left( \frac{\partial w^1}{\partial z} - \alpha_3 z T_1 \right) + c_{66} \left( \frac{1}{2} \left[ \frac{\partial u^1}{\partial y} + \frac{\partial v^1}{\partial x} \right] - \alpha_6 z T_1 \right). \end{aligned}$$

In these equations the  $c_{ij}$  are elastic constants generally called stiffnesses. To obtain  $c_{ij}$  from  $a_{ij}$  and vice versa, the following formulas [10] are to be used:

$$c_{ij} = \frac{\delta_{ij}}{\Delta},$$

and

$$a_{ij} = \frac{\delta'_{ij}}{\Delta'}.$$

Here  $\Delta$  and  $\Delta'$  stand for the determinants of the square matrices  $[a_{ij}]$  and  $[c_{ij}]$ , respectively;  $\delta_{ij}$  and  $\delta'_{ij}$  are the cofactors of  $a_{ij}$  and  $c_{ij}$ , respectively.

From the assumptions expressed by equations (2) and (3), the following relations result

$$\frac{\partial u^1}{\partial x} = -z \frac{\partial^2 w}{\partial x^2},$$

$$\frac{\partial v^1}{\partial y} = -z \frac{\partial^2 w}{\partial y^2},$$

$$\frac{1}{2} \left( \frac{\partial u^1}{\partial y} + \frac{\partial v^1}{\partial x} \right) = -z \frac{\partial^2 w}{\partial x \partial y} ,$$

and

$$\sigma_z^1 = 0 .$$

Note here that  $w$  is being written instead of  $w^1$ . This is due to the fact that  $T_0$  does not produce any bending, thus  $w$  and  $w^1$  are the same. Substituting these in the third equation of (12) yields

$$\begin{aligned} 0 = & c_{13} \left( -z \frac{\partial^2 w}{\partial x^2} - \alpha_1 z T_1 \right) + c_{23} \left( -z \frac{\partial^2 w}{\partial y^2} - \alpha_2 z T_1 \right) \\ & + c_{33} \left( \frac{\partial w}{\partial z} - \alpha_3 z T_1 \right) + c_{36} \left( -z \frac{\partial^2 w}{\partial x \partial y} - \alpha_6 z T_1 \right) . \end{aligned}$$

Thus

$$\begin{aligned} \frac{\partial w}{\partial z} - \alpha_3 z T_1 = & \frac{z}{c_{33}} \left[ c_{13} \left( \frac{\partial^2 w}{\partial x^2} + \alpha_1 T_1 \right) \right. \\ & \left. + c_{23} \left( \frac{\partial^2 w}{\partial y^2} + \alpha_2 T_1 \right) + c_{36} \left( \frac{\partial^2 w}{\partial x \partial y} + \alpha_6 T_1 \right) \right] . \end{aligned}$$

Substituting the expression for  $\left( \frac{\partial w}{\partial z} - \alpha_3 z T_1 \right)$  back into equations (12):

$$\begin{aligned} \sigma_x^1 = & z \left[ c_{11} \left( -\frac{\partial^2 w}{\partial x^2} - \alpha_1 T_1 \right) + c_{12} \left( -\frac{\partial^2 w}{\partial y^2} - \alpha_2 T_1 \right) \right. \\ & \left. + \frac{c_{13}}{c_{33}} \left\{ c_{13} \left( \frac{\partial^2 w}{\partial x^2} + \alpha_1 T_1 \right) + c_{23} \left( \frac{\partial^2 w}{\partial y^2} + \alpha_2 T_1 \right) + c_{36} \left( \frac{\partial^2 w}{\partial x \partial y} + \alpha_6 T_1 \right) \right\} \right] \\ = & z \left[ \left( \frac{c_{13} c_{13}}{c_{33}} - c_{11} \right) \frac{\partial^2 w}{\partial x^2} + \left( \frac{c_{23} c_{13}}{c_{33}} - c_{12} \right) \frac{\partial^2 w}{\partial y^2} + \left( \frac{c_{36} c_{13}}{c_{33}} - c_{36} \right) \frac{\partial^2 w}{\partial x \partial y} \right. \\ & \left. + \left\{ \alpha_1 \left( \frac{c_{13} c_{13}}{c_{33}} - c_{11} \right) + \alpha_2 \left( \frac{c_{23} c_{13}}{c_{33}} - c_{12} \right) + \alpha_6 \left( \frac{c_{36} c_{13}}{c_{33}} - c_{36} \right) \right\} T_1 \right] . \end{aligned}$$



The expression is simplified by defining

$$b_{ik} = - \left( \frac{c_{i3}c_{k3}}{c_{33}} - c_{ik} \right), \quad i, k = 1, 2, 6$$

and

$$\alpha_i^1 = b_{1i}\alpha_1 + b_{2i}\alpha_2 + b_{6i}\alpha_6 \quad i = 1, 2, 6$$

Following the same steps for the second and fourth equations of (12), the results obtained will be

$$\sigma_x^1 = -z \left( b_{11} \frac{\partial^2 w}{\partial x^2} + b_{12} \frac{\partial^2 w}{\partial y^2} + b_{16} \frac{\partial^2 w}{\partial x \partial y} + \alpha_1^1 T_1 \right), \quad (13)$$

$$\sigma_y^1 = -z \left( b_{12} \frac{\partial^2 w}{\partial x^2} + b_{22} \frac{\partial^2 w}{\partial y^2} + b_{26} \frac{\partial^2 w}{\partial x \partial y} + \alpha_2^1 T_1 \right),$$

and

$$\tau_{xy}^1 = -z \left( b_{16} \frac{\partial^2 w}{\partial x^2} + b_{26} \frac{\partial^2 w}{\partial y^2} + b_{66} \frac{\partial^2 w}{\partial x \partial y} + \alpha_6^1 T_1 \right).$$

The equations of equilibrium can be expressed as

$$\frac{\partial \sigma_x^1}{\partial x} + \frac{\partial \tau_{xy}^1}{\partial y} + \frac{\partial \tau_{xz}^1}{\partial z} = 0,$$

and

$$\frac{\partial \tau_{xy}^1}{\partial x} + \frac{\partial \sigma_y^1}{\partial y} + \frac{\partial \tau_{yz}^1}{\partial z} = 0. \quad (14)$$

From the first of these two equations

$$\begin{aligned} \frac{\partial \tau_{xz}^1}{\partial z} &= - \left( \frac{\partial \sigma_x^1}{\partial x} + \frac{\partial \tau_{xy}^1}{\partial y} \right) \\ &= z \left( b_{11} \frac{\partial^3 w}{\partial x^3} + b_{12} \frac{\partial^3 w}{\partial x \partial y^2} + b_{16} \frac{\partial^3 w}{\partial x^2 \partial y} + \alpha_1^1 \frac{\partial T_1}{\partial x} \right) \end{aligned}$$

$$+ z(b_{16} \frac{\partial^3 w}{\partial x^2 \partial y} + b_{26} \frac{\partial^3 w}{\partial y^3} + b_{66} \frac{\partial^3 w}{\partial x \partial y^2} + \alpha_6^1 \frac{\partial T_1}{\partial x}) ;$$

or simply

$$\frac{\partial \tau_{xz}^1}{\partial z} = z(A + B) , \quad (15)$$

where A and B are defined as

$$A = b_{11} \frac{\partial^3 w}{\partial x^3} + b_{12} \frac{\partial^3 w}{\partial x \partial y^2} + b_{16} \frac{\partial^3 w}{\partial x^2 \partial y} + \alpha_1^1 \frac{\partial T_1}{\partial x} ,$$

and

$$B = b_{16} \frac{\partial^3 w}{\partial x^2 \partial y} + b_{26} \frac{\partial^3 w}{\partial y^3} + b_{66} \frac{\partial^3 w}{\partial x \partial y^2} + \alpha_6^1 \frac{\partial T_1}{\partial y} .$$

Note here that

$$\tau_{yz}^1 \Big|_{z=\pm h} = \tau_{xz}^1 \Big|_{z=\pm h} = 0 .$$

If equation (15) is integrated with respect to z, then

$$\tau_{xz}^1 = \frac{z^2}{2} (A + B) + c .$$

At  $z = +h$  ,

$$0 = \frac{h^2}{2} (A + B) + c ,$$

or

$$c = - \frac{h^2}{2} (A + B) ;$$

and thus

$$\tau_{xz}^1 = \frac{z^2 - h^2}{2} (A + B) .$$

The shear force acting, per unit of length, across the thickness of the plate is

$$\begin{aligned}
 Q_x &= \int_{-h}^h \tau_{xz}^1 dz \\
 &= \int_{-h}^h \frac{z^2 - h^2}{2} (A + B) dz \\
 &= \frac{(A + B)}{2} \left[ \frac{z^3}{3} - h^2 z \right]_{-h}^h \\
 &= -\frac{2}{3} h^3 (A + B)
 \end{aligned}$$

Replacing A and B by their expressions in terms of w:

$$\begin{aligned}
 Q_x &= -\frac{2}{3} h^3 \left\{ b_{11} \frac{\partial^3 w}{\partial x^3} + b_{12} \frac{\partial^3 w}{\partial x \partial y^2} + b_{16} \frac{\partial^3 w}{\partial x^2 \partial y} + \alpha_1^1 \frac{\partial T_1}{\partial x} \right. \\
 &\quad \left. + b_{16} \frac{\partial^3 w}{\partial x^2 \partial y} + b_{26} \frac{\partial^3 w}{\partial y^3} + b_{66} \frac{\partial^3 w}{\partial x \partial y^2} + \alpha_6^1 \frac{\partial T_1}{\partial y} \right\} \quad (16-a)
 \end{aligned}$$

Following the same steps for the other equation of (14) yields

$$\begin{aligned}
 Q_y &= -\frac{2}{3} h^3 \left\{ b_{16} \frac{\partial^3 w}{\partial x^3} + b_{26} \frac{\partial^3 w}{\partial x \partial y^2} + b_{12} \frac{\partial^3 w}{\partial x^2 \partial y} + \alpha_6^1 \frac{\partial T_1}{\partial x} \right. \\
 &\quad \left. + b_{66} \frac{\partial^3 w}{\partial x^2 \partial y} + b_{22} \frac{\partial^3 w}{\partial y^3} + b_{26} \frac{\partial^3 w}{\partial x \partial y^2} + \alpha_2^1 \frac{\partial T_1}{\partial y} \right\} \quad (16-b)
 \end{aligned}$$

The condition of statical equilibrium [9] of an arbitrary volume element of the plate is represented by the equation

$$\frac{\partial Q_x}{\partial x} + \frac{\partial Q_y}{\partial y} + N_x \frac{\partial^2 w}{\partial x^2} + 2N_{xy} \frac{\partial^2 w}{\partial x \partial y} + N_y \frac{\partial^2 w}{\partial y^2} = 0 \quad .$$

Substituting the expressions for  $Q_x$  and  $Q_y$  yields

$$\begin{aligned} & - \frac{2}{3} h^3 \left\{ b_{11} \frac{\partial^4 w}{\partial x^4} + 2b_{16} \frac{\partial^4 w}{\partial x^3 \partial y} + (b_{12} + b_{66}) \frac{\partial^4 w}{\partial x^2 \partial y^2} + b_{26} \frac{\partial^4 w}{\partial x \partial y^3} \right. \\ & + \alpha_1^1 \frac{\partial^2 T_1}{\partial x^2} + \alpha_6^1 \frac{\partial^2 T_1}{\partial x \partial y} + b_{16} \frac{\partial^4 w}{\partial x^3 \partial y} + (b_{12} + b_{66}) \frac{\partial^4 w}{\partial x^2 \partial y^2} \\ & + \left. 2b_{26} \frac{\partial^4 w}{\partial x \partial y^3} + b_{22} \frac{\partial^4 w}{\partial y^4} + \alpha_6^1 \frac{\partial^2 T_1}{\partial x \partial y} + \alpha_2^1 \frac{\partial^2 T_1}{\partial y^2} \right\} \\ & + N_x \frac{\partial^2 w}{\partial x^2} + 2N_{xy} \frac{\partial^2 w}{\partial x \partial y} + N_y \frac{\partial^2 w}{\partial y^2} = 0 \quad . \end{aligned}$$

Rearranging of the above equation leads to a fourth order partial differential equation in  $w$

$$\begin{aligned} & b_{11} \frac{\partial^4 w}{\partial x^4} + 3b_{16} \frac{\partial^4 w}{\partial x^3 \partial y} + 2(b_{12} + b_{66}) \frac{\partial^4 w}{\partial x^2 \partial y^2} + 3b_{26} \frac{\partial^4 w}{\partial x \partial y^3} + b_{22} \frac{\partial^4 w}{\partial y^4} \\ & = - \left\{ \alpha_1^1 \frac{\partial^2 T_1}{\partial x^2} + 2\alpha_6^1 \frac{\partial^2 T_1}{\partial x \partial y} + \alpha_2^1 \frac{\partial^2 T_1}{\partial y^2} \right\} \\ & + \frac{3}{2h^3} \left\{ N_x \frac{\partial^2 w}{\partial x^2} + 2N_{xy} \frac{\partial^2 w}{\partial x \partial y} + N_y \frac{\partial^2 w}{\partial y^2} \right\} \quad . \end{aligned} \quad (17)$$

## CHAPTER III

### SOLUTION APPROACH

The thermoelastic problem described in the previous chapter is solved on a digital computer. The finite difference method is used to solve the equations (11) and (17) for the stress function  $F$  and deflection  $W$ , respectively. The computed values of  $F$  go into the equations (8) to furnish the stress tensor  $\tau^0$ . Also, the computed values of  $W$ , when substituted into the equations (13), give rise to stress tensor  $\tau^1$ . To help explain the method, consider a plate superimposed by a grid or lattice with a mesh size of  $\Delta h$  by  $\Delta k$  as shown in Figure 2.

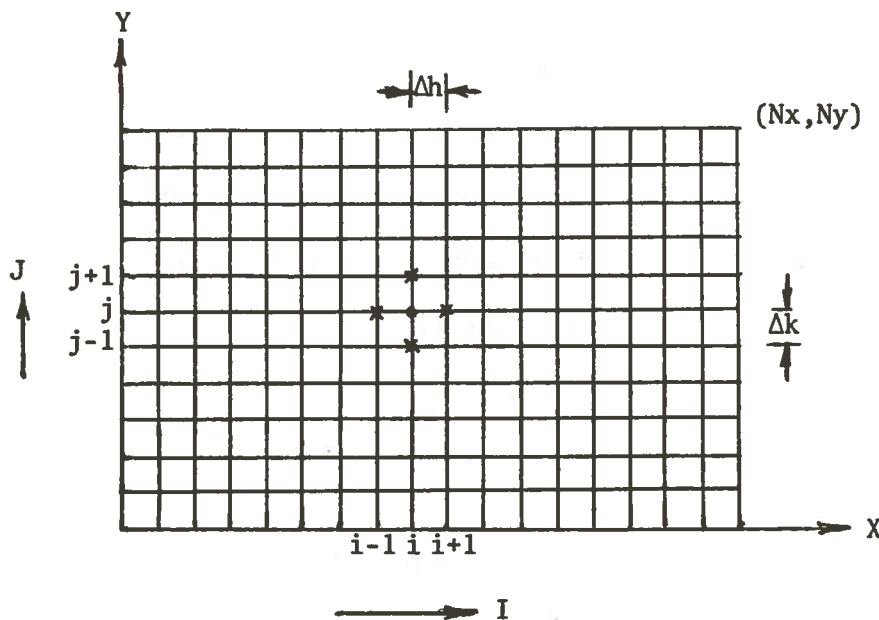


Figure 2

A Grid Superimposed on a Plate

In the case of equation (11) the numerical solution will consist of determining the values of the function  $F$  at the finitely spaced grid points shown. In this process, a combination of first, second, third, and fourth partial derivatives of  $F$  with respect to  $X$ ,  $Y$ , and  $X$  and  $Y$  will be required at the grid points. These partial derivatives are given in Table A-1 [11] of Appendix A. The following paragraphs briefly describe the method for obtaining them.

The Taylor series for the function  $F = F(x)$  at points  $(I+1)$  and  $(I-1)$  expanded about point  $(I)$  gives:

$$F_{i+1} = F_i + F'_i \Delta h + F''_i \frac{(\Delta h)^2}{2} + F'''_i \frac{(\Delta h)^3}{6} ,$$

and

(18)

$$F_{i-1} = F_i - F'_i \Delta h + F''_i \frac{(\Delta h)^2}{2} - F'''_i \frac{(\Delta h)^3}{6} ,$$

where  $F'_i$ ,  $F''_i$ , and  $F'''_i$  are first, second, and third partial derivatives of  $F_i$  with respect to  $X$ , respectively.

For obtaining the expressions for  $F'_i$  and  $F''_i$ , terms containing  $F'''_i$  and higher order partial derivatives may be neglected. Subtracting one equation from the other one obtains:

$$F_{i+1} - F_{i-1} = 2 F'_i \Delta h$$

or

$$\Delta h \frac{\partial F}{\partial x} = \Delta h F'_i = \frac{1}{2} \{ F_{i+1} - F_{i-1} \}$$

This is a central difference equation for the first partial derivative of  $F$  with respect to  $X$  using only three grid points.

Next, adding the two equations of (18) results in

$$F_{i-1} + F_{i+1} = 2 F_i + F_i'' (\Delta h)^2 ,$$

or

$$(\Delta h)^2 \frac{\partial^2 F}{\partial x^2} = (\Delta h)^2 F_i'' = F_{i-1} - 2F_i + F_{i+1}$$

This is a central difference equation for the second partial derivative of  $F$  with respect to  $X$  using only three grid points.

Similar expressions are obtained for third and fourth partial derivatives using more grid points.

The appropriate expressions for partial derivatives are substituted into the partial differential equation, and the resulting equation is solved at various grid points for the values of required function.

#### Finite Difference Equation for $F$

Referring to Table A-1, it may be noted that the partial derivatives required for solving equation (11) can be written as

$$\begin{aligned} \frac{\partial^4 F}{\partial x^4} &= \frac{1}{(\Delta h)^4} \{ F_{i-2,j} - 4F_{i-1,j} + 6F_{i,j} - 4F_{i+1,j} + F_{i+2,j} \} , \\ \frac{\partial^4 F}{\partial x^3 \partial y} &= \frac{\partial}{\partial y} \left( \frac{\partial^3 F}{\partial x^3} \right) \\ &= \frac{\partial}{\partial y} \left[ \frac{1}{2(\Delta h)^3} \{ -F_{i-2,j} + 2F_{i-1,j} - 2F_{i+1,j} + F_{i+2,j} \} \right] \\ &= \frac{1}{24(\Delta h)^3 \Delta k} [ -\{ F_{i-2,j-2} - 8F_{i-2,j-1} + 8F_{i-2,j+1} - F_{i-2,j+2} \} \\ &\quad + 2\{ F_{i-1,j-2} - 8F_{i-1,j-1} + 8F_{i-1,j+1} - F_{i-1,j+2} \} ] \end{aligned} \tag{19}$$

$$- 2\{F_{i+1,j-2} - 8F_{i+1,j-1} + 8F_{i+1,j+1} - F_{i+1,j+2}\} \\ + \{F_{i+2,j-2} - 8F_{i+2,j-1} + 8F_{i+2,j+1} - F_{i+2,j+2}\} ,$$

$$\frac{\partial^4 F}{\partial x^2 \partial y^2} = \frac{1}{144(\Delta h)^2 (\Delta k)^2} [-\{-F_{i-2,j-2} + 16F_{i-2,j-1} - 30F_{i-2,j} \\ + 16F_{i-2,j+1} - F_{i-2,j+2}\} + 16\{-F_{i-1,j-2} + 16F_{i-1,j-1} \\ - 30F_{i-1,j} + 16F_{i-1,j+1} - F_{i-1,j+2}\} - 30\{-F_{i,j-2} \\ + 16F_{i,j-1} - 30F_{i,j} + 16F_{i,j+1} - F_{i,j+2}\} \\ + 16\{-F_{i+1,j-2} + 16F_{i+1,j-1} - 30F_{i+1,j} + 16F_{i+1,j+1} \\ - F_{i+1,j+2}\} - \{-F_{i+2,j-2} + 16F_{i+2,j-1} - 30F_{i+2,j} \\ + 16F_{i+2,j+1} - F_{i+2,j+2}\}] , \quad (19) \text{ cont'd}$$

$$\frac{\partial^4 F}{\partial x \partial y^3} = \frac{1}{24\Delta h(\Delta k)^3} [ \{-F_{i-2,j-2} + 2F_{i-2,j-1} - 2F_{i-2,j+1} + F_{i-2,j+2}\} \\ - 8\{-F_{i-1,j-2} + 2F_{i-1,j-1} - 2F_{i-1,j+1} + F_{i-1,j+2}\} \\ + 8\{-F_{i+1,j-2} + 2F_{i+1,j-1} - 2F_{i+1,j+1} + F_{i+1,j+2}\} \\ - \{-F_{i+2,j-2} + 2F_{i+2,j-1} - 2F_{i+2,j+1} + F_{i+2,j+2}\}] ,$$

$$\frac{\partial^4 F}{\partial y^4} = \frac{1}{(\Delta k)^4} \{F_{i,j-2} - 4F_{i,j-1} + 6F_{i,j} - 4F_{i,j+1} + F_{i,j+2}\} .$$

The constants P, Q, R, S, and T are defined such that:

$$P = \frac{a_{22}}{(\Delta h)^4} , \quad Q = \frac{a_{26}}{12(\Delta h)^3 (\Delta k)} \\ R = \frac{(2a_{12} + a_{66})}{144(\Delta h)^2 (\Delta k)^2} , \quad S = \frac{a_{16}}{12\Delta h(\Delta k)^3} , \quad (20)$$



and

$$T = \frac{a_{11}}{(\Delta k)^4} \quad (20) \text{ cont'd}$$

Next, the right-hand side of equation (11) can be represented by a function RHSF (x,y) for a digital computer so that

$$\text{RHSF (x,y)} = \alpha_2 \frac{\partial^2 T_0}{\partial x^2} - 2\alpha_6 \frac{\partial^2 T_0}{\partial x \partial y} + \alpha_1 \frac{\partial^2 T_0}{\partial y^2} .$$

Substituting the expressions (19), the constants defined by equations (20), and the function RHSF (x,y) into the equation (11), and then rearranging the resulting equation, one would obtain the expression for  $F_{i,j}$ , the value of F at any point (I,J), as

$$\begin{aligned} F_{i,j} = & - [BM_1^F F_{i-2,j-2} + BM_2^F F_{i-2,j-1} + BM_3^F F_{i-2,j} + BM_4^F F_{i-2,j+1} \\ & + BM_5^F F_{i-2,j+2} + BM_6^F F_{i-1,j-2} + BM_7^F F_{i-1,j-1} + BM_8^F F_{i-1,j} \\ & + BM_9^F F_{i-1,j+1} + BM_{10}^F F_{i-1,j+2} + BM_{11}^F F_{i,j-2} + BM_{12}^F F_{i,j-1} \\ & + BM_{14}^F F_{i,j+1} + BM_{15}^F F_{i,j+2} + BM_{16}^F F_{i+1,j-2} + BM_{17}^F F_{i+1,j-1} \\ & + BM_{18}^F F_{i+1,j} + BM_{19}^F F_{i+1,j+1} + BM_{20}^F F_{i+1,j+2} + BM_{21}^F F_{i+2,j-2} \\ & + BM_{22}^F F_{i+2,j-1} + BM_{23}^F F_{i+2,j} + BM_{24}^F F_{i+2,j+1} + BM_{25}^F F_{i+2,j+2} \\ & + \text{RHSF (x,y)}] / BM_{13} . \end{aligned} \quad (21)$$

In the above equation, the coefficients  $BM_i$  are defined as

$$\begin{aligned}
BM_1 &= BM_{25} = Q + R + S \quad , \\
BM_2 &= BM_{24} = - 8Q - 16R - 2S \quad , \\
BM_3 &= BM_{23} = P + 30R \quad , \\
BM_4 &= BM_{22} = 8Q - 16R + 2S \quad , \\
BM_5 &= BM_{21} = - Q + R - S \quad , \\
BM_6 &= BM_{20} = - 2Q - 16R - 8S \quad , \\
BM_7 &= BM_{19} = 16Q + 256R + 16S \quad , \\
BM_8 &= BM_{18} = - 4P - 480R \quad , \\
BM_9 &= BM_{17} = - 16Q + 256R - 16S \quad , \\
BM_{10} &= BM_{16} = 2Q - 16R + 8S \quad , \\
BM_{11} &= BM_{15} = 30R + T \quad , \\
BM_{12} &= BM_{14} = - 480R - 4T \quad ,
\end{aligned}
\tag{22}$$

and

$$BM_{13} = 6P + 900R + 6T \quad .$$

### Finite Difference Equation for W

Due to the similarity between the left hand sides of the two equations (11) and (17), the left hand side of the finite difference form of equation (17) for W will be similar to the one for F, the only major change being in the coefficients.

Constants P1, Q1, R1, S1, and T1 are defined as

$$\begin{aligned}
P1 &= \frac{b_{11}}{(\Delta h)^4} \quad , & Q1 &= - \frac{3b_{16}}{24(\Delta h)^3 \Delta k} \quad , \\
R1 &= \frac{2(b_{12} + b_{66})}{144(\Delta h)^2 (\Delta k)^2} \quad , & S1 &= - \frac{3b_{26}}{24\Delta h (\Delta k)^3} \quad ,
\end{aligned}
\tag{23}$$

and

$$T1 = \frac{b_{22}}{(\Delta k)^4} \quad .$$

Thus the left hand side, LHS, of the equation (17) can be written as

$$\begin{aligned}
\text{LHS} = & DM_1 W_{i-2,j-2} + DM_2 W_{i-2,j-1} + DM_3 W_{i-2,j} + DM_4 W_{i-2,j+1} \\
& + DM_5 W_{i-2,j+2} + DM_6 W_{i-1,j-2} + DM_7 W_{i-1,j-1} + DM_8 W_{i-1,j} \\
& + DM_9 W_{i-1,j+1} + DM_{10} W_{i-1,j+2} + DM_{11} W_{i,j-2} + DM_{12} W_{i,j-1} \\
& + DM_{13} W_{i,j} + DM_{14} W_{i,j+1} + DM_{15} W_{i,j+2} + DM_{16} W_{i+1,j-2} \\
& + DM_{17} W_{i+1,j-1} + DM_{18} W_{i+1,j} + DM_{19} W_{i+1,j+1} + DM_{20} W_{i+1,j+2} \\
& + DM_{21} W_{i+2,j-2} + DM_{22} W_{i+2,j-1} + DM_{23} W_{i+2,j} + DM_{24} W_{i+2,j+1} \\
& + DM_{25} W_{i+2,j+2} ,
\end{aligned} \tag{24}$$

where the coefficients  $DM_i$  are expressed in terms of  $P_1$ ,  $Q_1$ ,  $R_1$ ,  $S_1$ , and  $T_1$  as

$$\begin{aligned}
DM_1 = DM_{25} &= Q_1 + R_1 + S_1 , \\
DM_2 = DM_{24} &= - 8Q_1 - 16R_1 - 2S_1 , \\
DM_3 = DM_{23} &= P_1 + 30R_1 , \\
DM_4 = DM_{22} &= 8Q_1 - 16R_1 + 2S_1 , \\
DM_5 = DM_{21} &= - Q_1 + R_1 - S_1 , \\
DM_6 = DM_{20} &= - 2Q_1 - 16R_1 - 8S_1 , \\
DM_7 = DM_{19} &= 16Q_1 + 256R_1 + 16S_1 , \\
DM_8 = DM_{18} &= - 4P_1 - 480R_1 , \\
DM_9 = DM_{17} &= - 16Q_1 + 256R_1 - 16S_1 , \\
DM_{10} = DM_{16} &= 2Q_1 - 16R_1 + 8S_1 , \\
DM_{11} = DM_{15} &= 30R_1 + T_1 , \\
DM_{12} = DM_{14} &= - 480R_1 - 4T_1 ,
\end{aligned} \tag{25}$$

and

$$DM_{13} = 6P1 + 900R1 + 6T1 \quad .$$

The right hand side of the equation (17) can be represented as

$$RHS = - RHSW(x,y) + \frac{3}{2h^3} \left\{ N_x \frac{\partial^2 w}{\partial x^2} + 2N_{xy} \frac{\partial^2 w}{\partial x \partial y} + N_y \frac{\partial^2 w}{\partial y^2} \right\} \quad , \quad (26)$$

where

$$RHSW(x,y) = \alpha_1 \frac{\partial^2 T_1}{\partial x^2} + 2\alpha_6 \frac{\partial^2 T_1}{\partial x \partial y} + \alpha_2 \frac{\partial^2 T_1}{\partial y^2} \quad .$$

Again, looking at Table A-1 in Appendix A one can write

$$\frac{\partial^2 w}{\partial x^2} = \frac{1}{(\Delta h)^2} \{ W_{i-1,j} - 2W_{i,j} + W_{i+1,j} \}$$

$$\frac{\partial^2 w}{\partial y^2} = \frac{1}{(\Delta k)^2} \{ W_{i,j-1} - 2W_{i,j} + W_{i,j+1} \} \quad (27)$$

$$\frac{\partial^2 w}{\partial x \partial y} = \frac{1}{4\Delta h \Delta k} \{ W_{i-1,j-1} - W_{i-1,j+1} - W_{i+1,j-1} + W_{i+1,j+1} \}$$

Substituting the expressions for LHS and RHS into the equation

(17) yields

$$\begin{aligned} W_{i,j} = & - [DM_1 W_{i-2,j-2} + DM_2 W_{i-2,j-1} + DM_3 W_{i-2,j} + DM_4 W_{i-2,j+1} \\ & + DM_5 W_{i-2,j+2} + DM_6 W_{i-1,j-2} + DM_7 W_{i-1,j-1} + \{DM_8 \\ & - h_1 N_{x_{i,j}} / (\Delta h)^2\} W_{i-1,j} + DM_9 W_{i-1,j+1} + DM_{10} W_{i-1,j+2} \\ & + DM_{11} W_{i,j-2} + \{DM_{12} - h_1 N_{y_{i,j}} / (\Delta k)^2\} W_{i,j-1} \\ & + \{DM_{14} - h_1 N_{y_{i,j}} / (\Delta k)^2\} W_{i,j+1} + DM_{15} W_{i,j+2} \end{aligned} \quad (28)$$

$$\begin{aligned}
& + DM_{16} W_{i+1,j-2} + DM_{17} W_{i+1,j-1} + \{DM_{18} - h_1 N_{x_{i,j}} / \\
& (\Delta h)^2\} W_{i+1,j} + DM_{19} W_{i+1,j+1} + DM_{20} W_{i+1,j+2} \\
& + DM_{21} W_{i+2,j-2} + DM_{22} W_{i+2,j-1} + DM_{23} W_{i+2,j} \\
& + DM_{24} W_{i+2,j+1} + DM_{25} W_{i+2,j+2} - h_1 N_{xy_{i,j}} \\
& \{W_{i-1,j-1} - W_{i-1,j+1} - W_{i+1,j-1} + W_{i+1,j+1}\} / \\
& (2\Delta h \Delta k) + RHSW(x,y) / [DM_{13} + 2h_1 \{N_{x_{i,j}} / (\Delta h)^2 \\
& + N_{y_{i,j}} / (\Delta k)^2\}] ,
\end{aligned}$$

(28) cont'd

where

$$h_1 = 3/2h^3 .$$

#### Boundary Conditions--All Edges Free

The problem described here is a boundary value problem; that is, the state of stress and/or deflection is defined at all four edges of the plate. The following are mathematical representations of the boundary conditions in terms of F and W.

#### A. Boundary Conditions in Terms of F

From the free edge condition, it follows that the normal forces are zero. Also, noting that there is no bending due to  $T_0$ , it is evident that the shear force is zero. Mathematically [2] the boundary conditions on the edge  $X = 0, a$  are

$$N_x = N_{xy} = 0 .$$

This leads to the relations [9]

$$F = \frac{\partial F}{\partial x} = 0 \quad ;$$

and on the edges  $Y = 0, b$  ,

$$F = \frac{\partial F}{\partial y} = 0 \quad .$$

### B. Boundary Conditions in Terms of W

At the plate's free edge the bending moment  $M_n$ , the shear force  $Q_n$ , and the twisting moment  $M_{ns}$  must all vanish, where  $(n,s)$  represents an orthogonal coordinate system,  $n$  being normal, and  $s$  being parallel to the edge. The twisting moment is, however, statically equivalent to a distributed shear force of intensity  $[-(\partial M_{ns}/\partial s)_{n=0}]$  .

Mathematically the boundary conditions are

$$\begin{aligned} M_n &= 0 \quad , \\ \text{and} \quad Q_n - \frac{\partial M_{ns}}{\partial s} &= 0 \end{aligned} \quad (29)$$

To obtain  $Q_n$  in terms of moments, consider the forces in the  $Z$  direction and moments about  $X$  and  $Y$  axis acting on an element of volume  $dx dy dz$  of the plate as shown in Figure 3. The moments  $M_x$ ,  $M_y$ , and  $M_{xy}$  are given as

$$\begin{aligned} M_x &= \int_{-h}^h \sigma_x z dz \quad , \\ \text{and} \quad M_y &= \int_{-h}^h \sigma_y z dz \quad , \\ M_{xy} &= \int_{-h}^h \tau_{xy} z dz \quad ; \end{aligned} \quad (30)$$

and the shear forces  $Q_x$  and  $Q_y$  are represented by the equations (16-A) and (16-B).

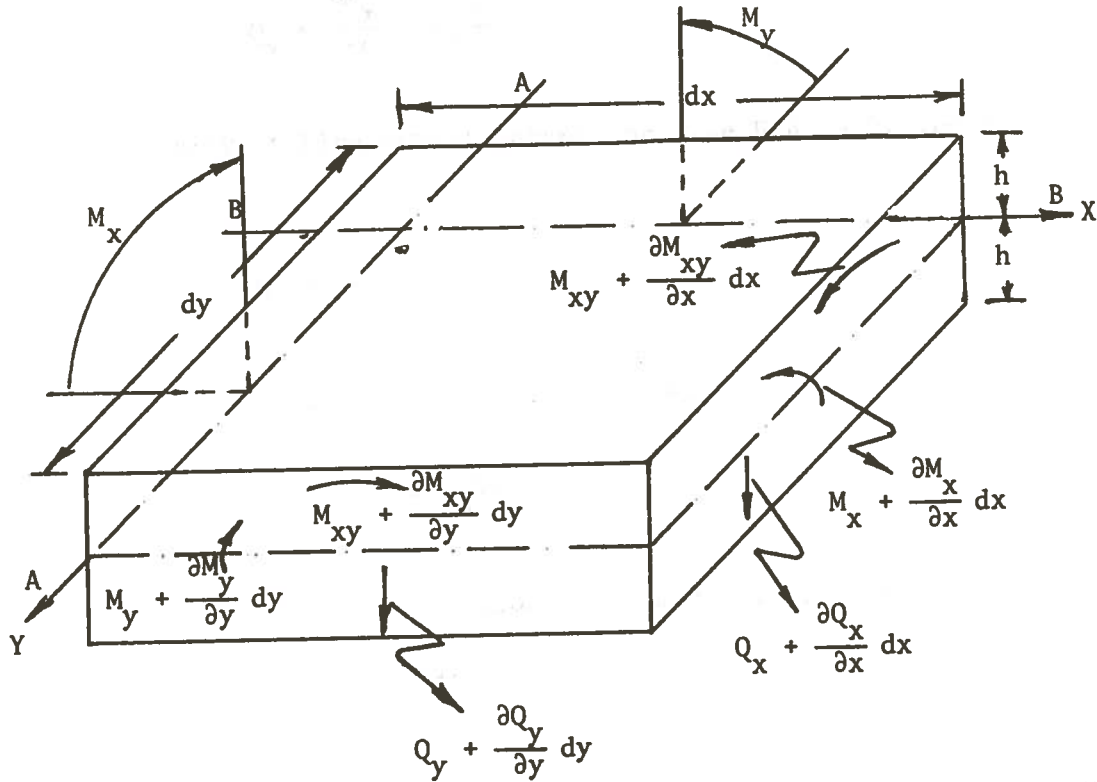


Figure 3

Moments and Shear Forces Acting on a Volume Element

Taking moments about the line A-A in Figure 3 one obtains

$$\begin{aligned} & (M_x + \frac{\partial M_x}{\partial x} dx) dy - M_x dy + M_{xy} dx - (M_{xy} + \frac{\partial M_{xy}}{\partial y} dy) dx \\ & - (Q_x + \frac{\partial Q_x}{\partial x} dx) dy dx + Q_y dx \frac{dx}{2} - (Q_y + \frac{\partial Q_y}{\partial y} dy) dx \frac{dx}{2} = 0 \end{aligned}$$

By neglecting the terms containing the product  $dx dy dz$ , the above equation becomes

$$Q_x = \frac{\partial M_x}{\partial x} - \frac{\partial M_{xy}}{\partial y} .$$

Similarly, taking moments about the line B-B in Figure 3 results in

$$Q_y = \frac{\partial M_y}{\partial y} - \frac{\partial M_{xy}}{\partial x} .$$

Replacing  $x$  and  $y$  by  $n$  and  $s$ , the shear force equation becomes

$$Q_n = \frac{\partial M_n}{\partial n} - \frac{\partial M_{ns}}{\partial s} . \quad (31)$$

Interaction of the equations (29) and (31) results in the following boundary conditions:

$$M_n = 0 ,$$

and

$$\frac{\partial M_n}{\partial n} - 2 \frac{\partial M_{ns}}{\partial s} = 0 .$$

(32)

On the Edges  $X = 0, a$

From the equations (13) and (30)

$$\begin{aligned} M_x &= - \int_{-h}^h z^2 (b_{11} \frac{\partial^2 w}{\partial x^2} + b_{12} \frac{\partial^2 w}{\partial y^2} + b_{16} \frac{\partial^2 w}{\partial x \partial y} + \alpha_1^1 T_1) \\ &= - \frac{2}{3} h^3 (b_{11} \frac{\partial^2 w}{\partial x^2} + b_{12} \frac{\partial^2 w}{\partial y^2} + b_{16} \frac{\partial^2 w}{\partial x \partial y} + \alpha_1^1 T_1) . \end{aligned}$$



Similarly,  $M_y$  and  $M_{xy}$  are

$$M_y = -\frac{2}{3} h^3 (b_{12} \frac{\partial^2 w}{\partial x^2} + b_{22} \frac{\partial^2 w}{\partial y^2} + b_{26} \frac{\partial^2 w}{\partial x \partial y} + \alpha_2^1 T_1) ,$$

and

$$M_{xy} = -\frac{2}{3} h^3 (b_{16} \frac{\partial^2 w}{\partial x^2} + b_{26} \frac{\partial^2 w}{\partial y^2} + b_{66} \frac{\partial^2 w}{\partial x \partial y} + \alpha_6^1 T_1) .$$

Further, differentiating  $M_x$  and  $M_{xy}$ :

$$\begin{aligned} \frac{\partial M_x}{\partial x} - 2 \frac{\partial M_{xy}}{\partial y} &= -\frac{2}{3} h^3 (b_{11} \frac{\partial^3 w}{\partial x^3} + b_{12} \frac{\partial^3 w}{\partial x \partial y^2} + b_{16} \frac{\partial^3 w}{\partial x^2 \partial y} \\ &\quad + \alpha_1^1 \frac{\partial T_1}{\partial x}) + \frac{4}{3} h^3 (b_{16} \frac{\partial^3 w}{\partial y \partial x^2} + b_{26} \frac{\partial^3 w}{\partial y^3} \\ &\quad + b_{66} \frac{\partial^3 w}{\partial x \partial y^2} + \alpha_6^1 \frac{\partial T_1}{\partial y}) \\ &= -\frac{2}{3} h^3 (b_{11} \frac{\partial^3 w}{\partial x^3} - b_{16} \frac{\partial^3 w}{\partial x^2 \partial y} + (b_{12} - 2b_{66}) \frac{\partial^3 w}{\partial x \partial y^2} \\ &\quad - 2b_{26} \frac{\partial^3 w}{\partial y^3} + \alpha_1^1 \frac{\partial T_1}{\partial x} - 2\alpha_6^1 \frac{\partial T_1}{\partial y}) . \end{aligned}$$

Thus the equations (32) reduce to

$$b_{11} \frac{\partial^2 w}{\partial x^2} + b_{12} \frac{\partial^2 w}{\partial y^2} + b_{16} \frac{\partial^2 w}{\partial x \partial y} + \alpha_1^1 T_1 = 0 ,$$

and

$$\begin{aligned} b_{11} \frac{\partial^3 w}{\partial x^3} - b_{16} \frac{\partial^3 w}{\partial x^2 \partial y} + (b_{12} - 2b_{66}) \frac{\partial^3 w}{\partial x \partial y^2} - 2b_{26} \frac{\partial^3 w}{\partial y^3} \\ + \alpha_1^1 \frac{\partial T_1}{\partial x} - 2\alpha_6^1 \frac{\partial T_1}{\partial y} = 0 . \end{aligned} \quad (33)$$

On the Edges  $Y = 0, b$

Proceeding similarly to obtain the boundary conditions in terms of  $W$  for the edges normal to  $Y$ -axis, the following equations are to be satisfied on the edges:

$$b_{12} \frac{\partial^2 w}{\partial x^2} + b_{22} \frac{\partial^2 w}{\partial y^2} + b_{26} \frac{\partial^2 w}{\partial x \partial y} + \alpha_2^1 T_1 = 0 \quad ,$$

and

$$2b_{16} \frac{\partial^3 w}{\partial x^3} + (2b_{66} - b_{12}) \frac{\partial^3 w}{\partial x^2 \partial y} + b_{26} \frac{\partial^3 w}{\partial x \partial y^2} \quad (34)$$

$$- b_{22} \frac{\partial^3 w}{\partial y^3} + 2\alpha_6^1 \frac{\partial T_1}{\partial x} - \alpha_2^1 \frac{\partial T_1}{\partial y} = 0 \quad .$$

## CHAPTER IV

### THE COMPUTER PROGRAM

The computer program is written in Fortran IV language for the IBM System 360, Model 65 computer. It is designed to compute the thermal stresses and deflections in an anisotropic plate, having at least one plane of elastic symmetry, being free of external forces and having all four edges free. A flow chart for the main program is shown in Figure 4. The flow charts for the subroutines written specifically to solve the stress analysis problem are shown in Figures 5 to 8 inclusive. The temperature distributions that are to be read into the computer are represented by fourth order polynomials in X, and Y, namely,

$$\begin{aligned} T_0 = & u_1x^4 + u_2x^3y + u_3x^3 + u_4x^2y^2 + u_5x^2y + u_6x^2 + u_7x \\ & + u_8xy + u_9y + u_{10}y^2 + u_{11}xy^2 + u_{12}y^3 + u_{13}xy^3 \\ & + u_{14}y^4 + u_{15}, \end{aligned}$$

(35)

and

$$\begin{aligned} T_1 = & v_1x^4 + v_2x^3y + v_3x^3 + v_4x^2y^2 + v_5x^2y + v_6x^2 + v_7x \\ & + v_8xy + v_9y + v_{10}y^2 + v_{11}xy^2 + v_{12}y^3 + v_{13}xy^3 \\ & + v_{14}y^4 + v_{15}, \end{aligned}$$

where  $u_1, u_2, \dots, u_{15}$ , and  $v_1, v_2, \dots, v_{15}$  are coefficients of the polynomials.

Data is read into the computer in the following sequence:

CARD 1

$N_x, N_y$  (the number of grid points along  $x$ , and  $y$  directions of the plate, respectively),  $\Delta h, \Delta k$  (the increments along  $x$ , and  $y$  directions, respectively), and  $h$ , the half thickness of the plate.

CARDS 2 to 7

Elastic constants or compliances  $a_{ij}$  (there will be 36 of them).

CARD 8

The linear coefficients of thermal expansion in the  $i$ th direction,  $\alpha_i$ .

CARDS 9 to 12

The coefficients for the polynomials given by (35). These are  $u_1, u_2, \dots, u_{15}$ , and  $v_1, v_2, \dots, v_{15}$ .

CARDS 13 to 16

Repeat coefficients from cards 9 to 12. This is done to facilitate data transfer within the program.

The entire program with sample input is listed in the Appendix A.

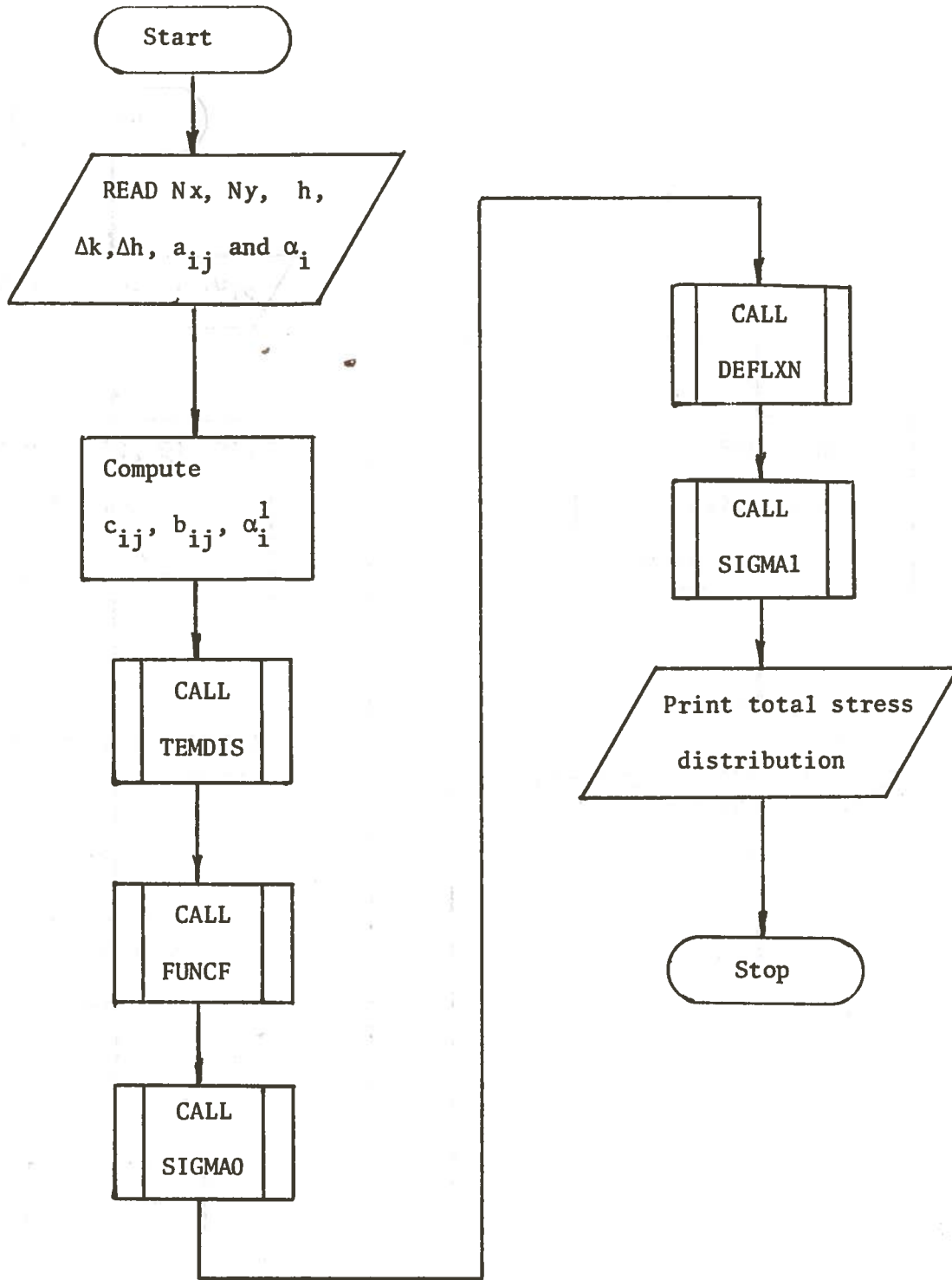


Figure 4  
Main Program Flow Chart

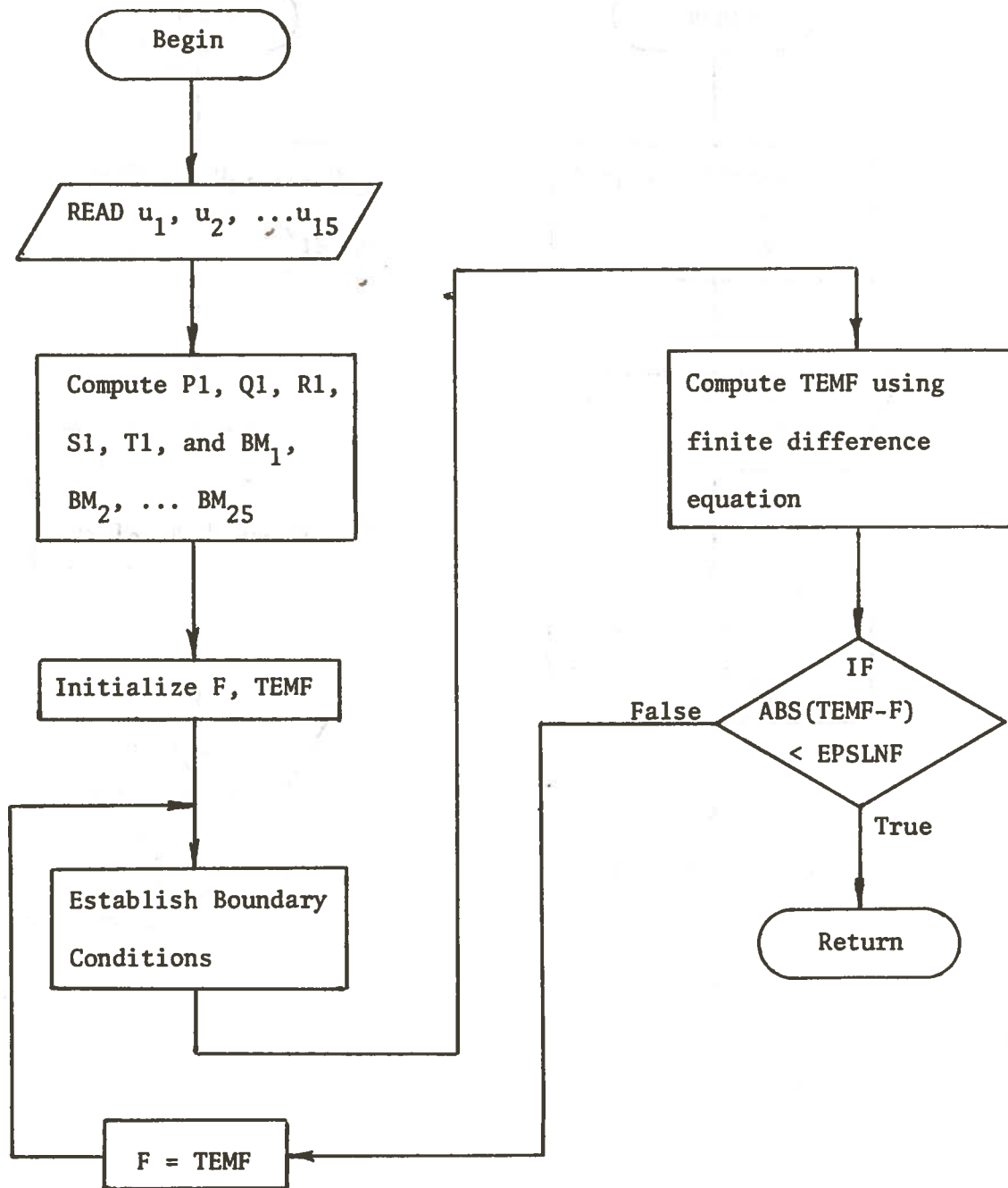


Figure 5

Subroutine FUNCF Flow Chart

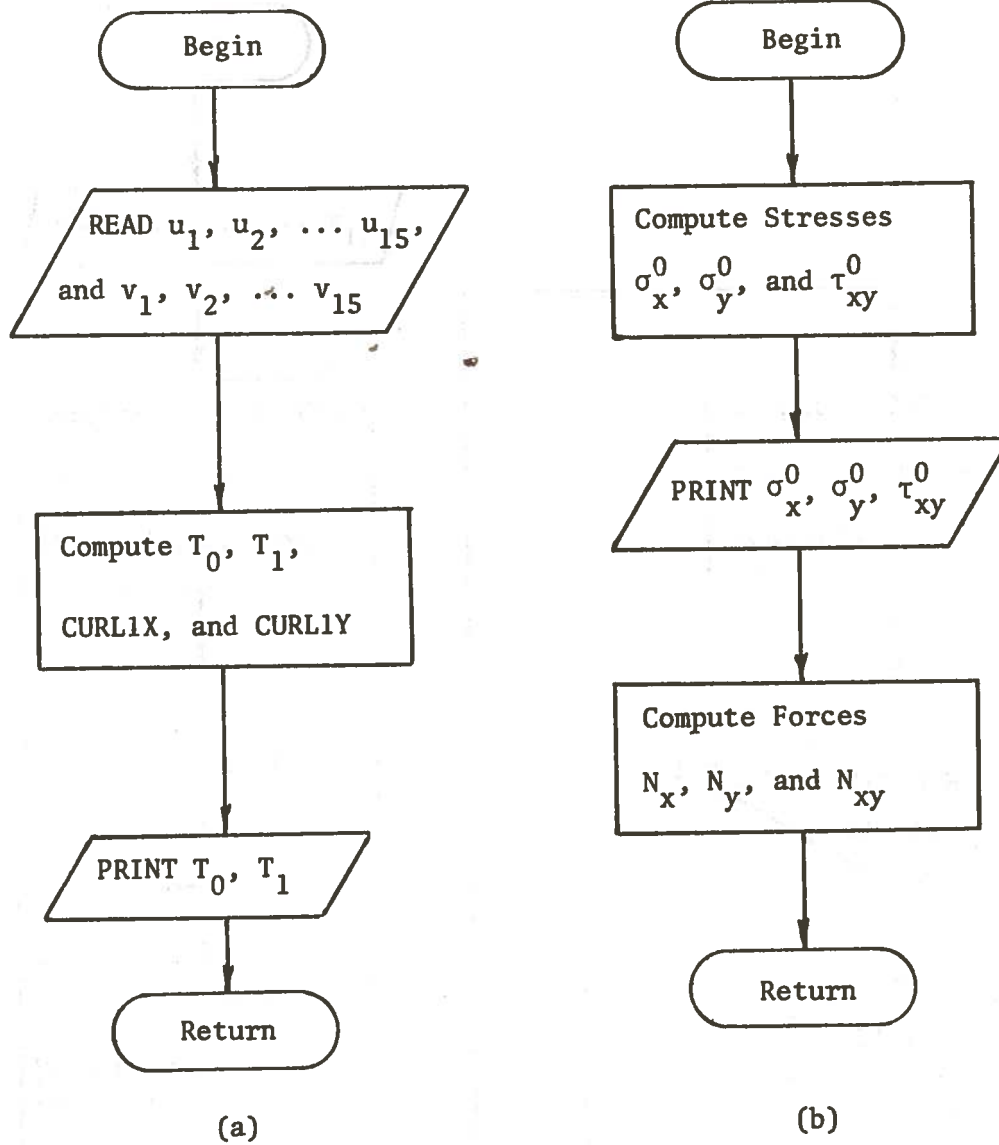


Figure 6

Flow Charts for Subroutines TEMDIS (6a) and SIGMA0 (6b)

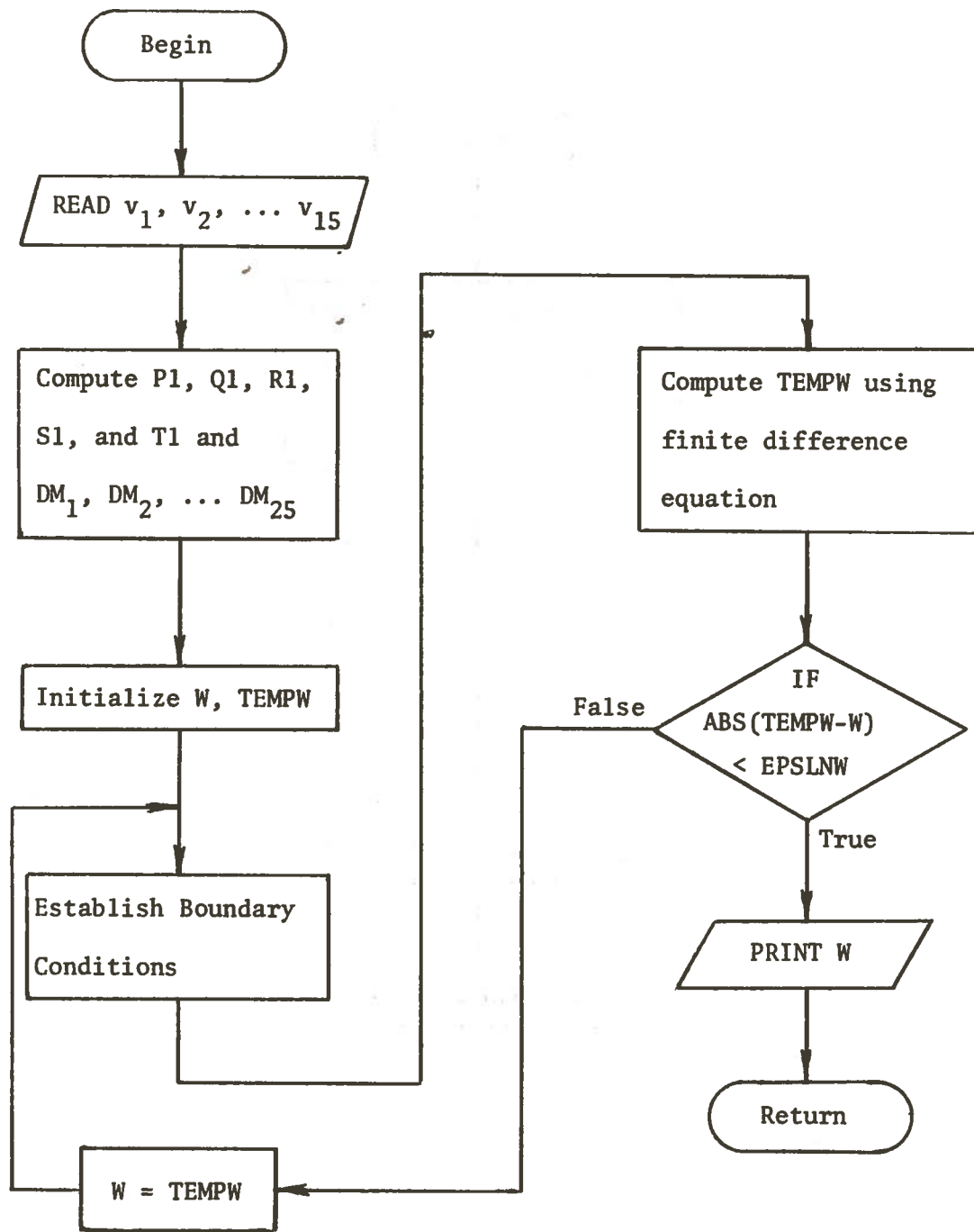


Figure 7

Subroutine DEFLXN Flow Chart



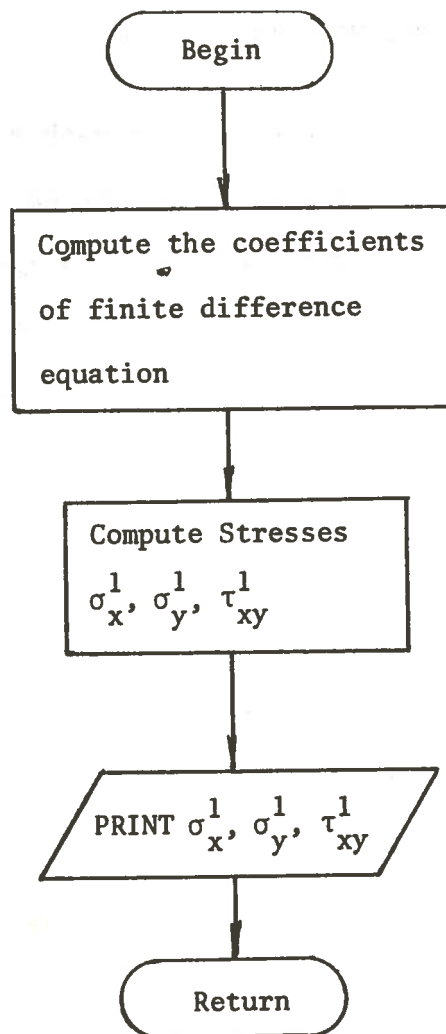


Figure 8

Subroutine SIGMA1 Flow Chart

## RESULTS AND CONCLUSIONS

The program developed is capable of computing thermal stresses and deflections in an anisotropic plate having one plane of elastic symmetry and no edge constraints. Results for a sample problem (input-data included in Appendix A) are given in Appendix B. The results are not completely satisfactory. This is due to the fact that the elastic constants,  $a_{ij}$ , are not exact. These constants are not available in the literature. These are known for a number of crystalline materials but, unfortunately, they are not of interest. Thus the elastic constants were assumed to be similar to those of steel which is an isotropic material. The program was run on a trial and error basis, changing some of these elastic constants on every run. The results given in Appendix B seem to be reasonable for structural materials.

It is suggested that further work in this field be directed toward experimentally obtaining the elastic constants of interest, so that they may be analyzed for stresses and deflections due to temperature gradients. The program itself could be modified by attempting to solve the set of linear finite difference equations implicitly.

TABLE A-1  
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APPENDIX A

TABLE A-1

$h^n f(n)$	Multiplier	Differentiation Formula Coefficients				
$hf'$	1	(-1)	1			
	1/2	(-3)	4	-1		
	1/6	(-11)	18	-9	2	
	1/12	(-25)	48	-36	16	-3
	1/2	-1	(0)	1		
	1/6	-2	(-3)	6	-1	
	1/12	-3	(-10)	18	-6	1
	1/12	1	-8	(0)	8	-1
$h^2 f''$	1	(1)	-2	1		
	1	(2)	-5	4	-1	
	1/12	(35)	-104	114	-56	11
	1	1	(-2)	1		
	1	1	(-2)	1	0	
	1/12	11	(-20)	6	4	-1
	1/12	-1	16	(-30)	16	-1
$h^3 f'''$	1	(-1)	3	-3	1	
	1/2	(-5)	18	-24	14	-3
	1	-1	(3)	-3	1	
	1/2	-3	(10)	-12	6	-1
	1/2	-1	2	(0)	-2	1
$h^4 f^{IV}$	1	(1)	-4	6	-4	1
	1	1	(-4)	6	-4	1
	1	1	-4	(6)	-4	1

The number within the parentheses is the coefficient of the (i,j)th term of the finite difference expression. Those to its right are the coefficients of the finite difference terms in the forward direction, and the others are the coefficients of the finite difference terms in the backward direction.

## Computer Program Listing

```

C THE PROGRAM IS DESIGNED TO COMPUTE THERMAL STRESSES
C AND DEFLECTIONS IN AN ANISOTROPIC PLATE. THE PLATE
C DIMENSIONS ARE REPRESENTED BY NX, NY, AND DELH, DELK.
C ADJUSTING THESE PARAMETERS DIFFERENT SIZES OF PLATE
C CAN BE ANALYZED FOR THERMAL STRESSES. IF NEED BE, THE
C DIMENSIONS OF F, TEMF, ETC., IN THE DIMENSION STATEMENT
C MAY ALSO BE CHANGED.
C
C DIMENSION F(30,30),TEMF(30,30),ANX(30,30),ANY(30,30)
C DIMENSION ANXY(30,30),W(30,30),TEMPW(30,30)
C DIMENSION TELX(30,30),TELY(30,30),TELXY(30,30)
C DIMENSION T0(30,30),T1(30,30),CURLIX(30,30),CURLIY(30,30)
C DIMENSION A(6,6),AT(6,6),C(6,6),B(6,6)
C DIMENSION ALFA(6),K(6),G(6)
C INTEGER G
C
C READ IN COMPLIANCES, DIMENSIONS OF PLATE, INCREMENTS
C ALONG X AND Y DIRECTIONS, AND LINEAR COEFFICIENTS OF
C THERMAL EXPANSION.
C READ(5,15) NX,NY,DELH,DELK,H
C
C DO 700 I=1,6
C 700 READ(5,710) (A(I,J),J=1,6)
C DO 6 I=1,6
C 6 WRITE (6,8) (A(I,J),J=1,6)
C 8 FORMAT(10X,6F15,4)
C READ(5,10) AA1,AA2,AA3,AA6
C N=6
C
C A1=AA1#1.E-6
C A2=AA2#1.E-6
C A3=AA3#1.E-6

```

```

C
C
C
A6=AA6*1.E-6
C
C
C
COMPUTE THE ELASTIC CONSTANTS C(I,J).
C
C
C
DO 600 I=1,6
DO 600 J=1,6
AT(I,J)=A(I,J)
CALL MINV(AT,N,D,K,G)
C
C
C
MINV IS A SUBROUTINE AVAILABLE IN THE SYSLIB.
THIS SUBROUTINE GIVES THE VALUES OF DETERMINANT
AND THE COFACTORS OF THE ELEMENTS A(I,J)S OF THE MATRIX A.
C
C
C
DO 610 I=1,6
DO 610 J=1,6
DO 620 L=1,6
DO 620 M=1,6
620 AT(L,M)=A(L,M)
DO 630 M=1,6
IF(M.EQ.J) GO TO 640
AT(I,M)=0.
GO TO 630
640 AT(I,M)=1.
630 CONTINUE
CALL MINV(AT,N,D1,K,G)
L=I+J
DO 650 M=1,L
D1=D1*(-1)
650 C(I,J)=D1/D
610 CONTINUE
C
C
C
DETERMINE THE CONSTANTS B(I,J) AND MODIFIED THERMAL
EXPANSION COEFFICIENTS ALFAS.

```

```

C
DO 720 I=1,6
IF(I.GT.2.AND.I.LT.6) GO TO 720
DO 730 J=1,6
IF(J.GT.2.AND.J.LT.6) GO TO 730
B(I,J)=(C(I,J)-C(I,3))*(C(J,3)/C(3,3))*1.E+8
CONTINUE
730 ALFA(I)=B(1,I)*A1+B(2,I)*A2+B(6,I)*A6
CONTINUE
720 H1=3./(2*H**3.)
ALFA1=ALFA(1)
ALFA2=ALFA(2)
ALFA6=ALFA(6)
C
CALL TEMDIS(DELH,DELK,NX,NY,T0,T1,CURLIX,CURLY,H)
C
CALL FUNCF(A,A1,A2,A3,A6,NX,NY,DELH,DELK,F,TEMP)
C
CALL SIGMA0 (F,NX,NY,H,DELH,DELK,ANX,ANY,ANXY)
C
CALL DEFLXN (B,ANX,ANY,ANXY,NX,NY,DELH,DELK,W,H1,
• ALFA1,ALFA2,ALFA3,ALFA6,TEMPW,T1,CURLIX,CURLY,ALFA)
C
CALL SIGMA1(W,B,DELH,DELK,H,NX,NY,TELEX,TELY,TELXY,ALFA,T1)
C
C COMPUTE TOTAL STRESS DISTRIBUTION ON THE UPPER SURFACE
DO 900 I=1,NX
DO 900 J=1,NY
ANX(I,J)=ANX(I,J)+TELEX(I,J)
ANY(I,J)=ANY(I,J)+TELY(I,J)
900 ANXY(I,J)=ANXY(I,J)+TELXY(I,J)
C

```

```

C      PRINT TOTAL STRESS DISTRIBUTION ON THE UPPER SURFACE
      WRITE (6,155) (J,J=1,6)
      DO 160 I=1,NX
160    WRITE (6,210) I,(ANX(I,J),J=1,NY)
      WRITE (6,165) (J,J=1,6)
      DO 172 I=1,NX
172    WRITE (6,210) I,(ANY(I,J),J=1,NY)
      WRITE (6,175) (J,J=1,6)
      DO 180 I=1,NX
180    WRITE (6,210) I,(ANXY(I,J),J=1,NY)
      FORMAT(4F10.4)
15    FORMAT(2I6,3F8.3)
155    FORMAT(1H1//////////1X,T30,'STRESS SIGMAX ON THE UPPER SURFACE (P
      •SI)///17X,6(12,12X)//)
165    FORMAT(1H1//////////1X,T30,'STRESS SIGMAY ON THE UPPER SURFACE (P
      •SI)///17X,6(12,12X)//)
175    FORMAT(1H1//////////1X,T30,'STRESS SIGMAXY ON THE UPPER SURFACE (
      •PSI)///17X,6(12,12X)//)
210    FORMAT(1H,I12,(T15,6E14.5))
      710 FORMAT(6F10.4)
      STOP
      END

```

```

SUBROUTINE FUNCF(A,A1,A2,A3,A6,NX,NY,DELH,DELK,F,TEMP)
DIMENSION F(30,30),TEMP(30,30),A(6,6),ANXY(30,30)
DIMENSION T0(30,30),T1(30,30),RHS(30,30)
X(Z)=(Z-1.)*DELH
Y(Z)=(Z-1.)*DELK

```

```

C
C      RIGHT HAND SIDE FUNCTION RHSF(X,Y)

```



```

RHSF(X,Y)= ((A2*(12.*U1*X*X+6.*(U2*X*Y+U3*X)+2.*(U4*Y*Y
+U5*Y+U6))+A1*(2.*U4*X*X+2.*U10+2.*U11*X+6.*U12*Y
+6.*U13*X*Y+12.*U14*Y*Y)-2.*A6*(3.*U2*X*X+4.*U4*X*Y
+2.*U5*X+U8+2.*U11*Y+3.*U13*Y*Y))*1.E+8)
EPSLNF=1.E-2
READ(5,10) U1,U2,U3,U4,U5,U6,U7,U8,U9,U10,U11,U12,U13,
C      U14
NX2=NX-2
NY2=NY-2
C
C      COMPUTE THE CONSTANTS P,Q,R,S,T, AND BMS.
C
P=A(2,2)/(DELH**4)
Q=A(2,6)/(12.*DELH**3*DELK)
R= (2.*A(1,2)+A(6,6))/(144.*DELH**2*DELK**2)
S=A(1,6)/(12.*DELH*DELK**3)
T=A(1,1)/(DELK**4)
C
BM1=Q+R+S
BM2=-8.*Q-16.*R-2.*S
BM3=P+30.*R
BM4=8.*Q-16.*R+2.*S
BM5=-Q+R-S
BM6=-2.*Q-16.*R-8.*S
BM7=16.*Q+256.*R+16.*S
BM8=-4.*P-480.*R
BM9=-16.*Q+256.*R-16.*S
BM10=2.*Q-16.*R+8.*S
BM11=30.*R+T
BM12=-480.*R-4.*T
BM13=6.*P+900.*R+6.*T
BM14=BM12

```

```

BM15=BM11
BM16=BM10
BM17=BM9
BM18=BM8
BM19=BM7
BM20=BM6
BM21=BM5
BM22=BM4
BM23=BM3
BM24=BM2
BM25=BM1

C
C INITIALIZE F, AND TEMF AND ESTABLISH THE BOUNDARY
C CONDITIONS.
C
DO 1500 I=1,NX
DO 1500 J=1,NY
F(I,J)=0.
1500 TEMF(I,J)=0.
M=1
100 DO 90 I=3,NX2
DO 90 J=3,NY2
U=I
V=J
TEMF(I,J)=-((BM1*F(I-2,J-2)+BM2*F(I-2,J-1)+BM3*F(I-2,J))+
BM4*F(I-2,J+1)+BM5*F(I-2,J+2)+BM6*F(I-1,J-2)+BM7*F(I-1,J-1)
+BM8*F(I-1,J)+BM9*F(I-1,J+1)+BM10*F(I-1,J+2)+BM11*F(I,J-2)
+BM12*F(I,J-1)+BM14*F(I,J+1)+BM15*F(I,J+2)+BM16*F(I+1,J-2)
+BM17*F(I+1,J-1)+BM18*F(I+1,J)+BM19*F(I+1,J+1)+BM20*
F(I+1,J+2)+BM21*F(I+2,J-2)+BM22*F(I+2,J-1)+BM23*F(I+2,J)+
BM24*F(I+2,J+1)+BM25*F(I+2,J+2))+RHSF(X(U),Y(V)))/BM13
90 CONTINUE

```

```

C
C CHECK FOR CONVERGENCE OF F.
DO 120 I=1,NX
DO 120 J=1,NY
IF (ABS(TEMP(I,J)-F(I,J)).GT.EPSLNF) GO TO 140
CONTINUE
GO TO 150
140 M=M+1
DO 1010 I=1,NX
DO 1010 J=1,NY
1010 F(I,J)=TEMP(I,J)
IF (M.LE.40) GO TO 100
10 FORMAT(8F10.4)
150 RETURN
END

```

```

SUBROUTINE SIGMA0 (F,NX,NY,H,DELH,DELK,ANX,ANY,ANXY)
DIMENSION F(30,30),ANX(30,30),ANY(30,30),ANXY(30,30)
C ANX, ANY, ANXY RESPECTIVELY REPRESENT STRESSES SIGMAX-0,
C SIGMAY-0, AND SIGMAXY-0.
NXA=NX-4
DO 120 I=1,NXA
DO 120 J=1,4
ANX(I,J)=(F(I,J)-2.*F(I,J+1)+F(I,J+2))/DELK**2.
ANY(I,J)=(F(I,J)-2.*F(I+1,J)+F(I+2,J))/DELH**2.
ANXY(I,J)=(F(I,J)-F(I,J+1)-F(I+1,J)+F(I+1,J+1))/(DELH*DELK)
DO 130 I=1,4
DO 130 J=5,NY
ANX(I,J)=(F(I,J)-2.*F(I,J-1)+F(I,J-2))/DELK**2.

```

```

130 ANY(I,J)=(F(I,J)-2.*F(I+1,J)+F(I+2,J))/DELH**2.
    ANXY(I,J)=-F(I,J)-F(I,J-1)-F(I+1,J)+F(I+1,J-1)/(DELH*DELK)
    NYA=NY-3
    DO 140 I=5,NX
    DO 140 J=NYA ,NY
    ANX(I,J)=(F(I,J)-2.*F(I,J-1)+F(I,J-2))/DELK**2.
    ANY(I,J)=(F(I,J)-2.*F(I-1,J)+F(I-2,J))/DELH**2.
140 ANXY(I,J)=(F(I,J)-F(I,J-1)-F(I-1,J)+F(I-1,J-1))/(DELH*DELK)
    NXA=NX-3
    NYA=NY-4
    DO 145 I=NXA ,NX
    DO 145 J=1,NYA
    ANX(I,J)=(F(I,J)-2.*F(I,J+1)+F(I,J+2))/DELK**2.
    ANY(I,J)=(F(I,J)-2.*F(I-1,J)+F(I-2,J))/DELH**2.
145 ANXY(I,J)=-F(I,J)-F(I,J+1)-F(I-1,J)+F(I-1,J+1)/(DELH*DELK)
    NXA=NX-4
    DO 150 I= 5,NXA
    DO 150 J= 5,NYA
    ANX(I,J)=2*H*(-F(I,J-2)+16.*F(I,J-1)-30.*F(I,J)+16.*
    *F(I,J+1)-F(I,J+2))/(12.*DELK**2.)
    ANY(I,J)=2*H*(-F(I-2,J)+16.*F(I-1,J)-30.*F(I,J)+16.*
    *F(I+1,J)-F(I+2,J))/(12.*DELH**2.)
150 ANXY(I,J)=2*H*(F(I-1,J-1)-F(I-1,J+1)-F(I+1,J-1)
    *+F(I+1,J+1))/(4.*DELH*DELK)
C
C PRINT THE STRESSES SIGMAX-0, SIGMAY-0, SIGMAXY-0.
WRITE (6,155) (J,J=1,6)
DO 160 I=1,NX
160 WRITE (6,210) I,(ANX(I,J),J=1,NY)
    WRITE (6,165) (J,J=1,6)
DO 172 I=1,NX
172 WRITE (6,210) I,(ANY(I,J),J=1,NY)

```

```

WRITE (6,175) (J,J=1,6)
DO 180 I=1,NX
180 WRITE (6,210) I,(ANXY(I,J),J=1,NY)
C
C ANX, ANY, ANXY RESPECTIVELY REPRESENT FORCES NX, NY, NXY.
C FROM HERE ON.
DO 170 I=1,NX
DO 170 J=1,NY
ANX(I,J)=ANX(I,J)*2.#H
ANY(I,J)=ANY(I,J)*2.#H
170 ANXY(I,J)=ANXY(I,J)*2.#H
155 FORMAT(1H1//////////1X,T40,'STRESS SIGMAX-0 (PSI)')////
. 17X,6(I2,12X)//
165 FORMAT(1H1//////////1X,T40,'STRESS SIGMAY-0 (PSI)')////
. 17X,6(I2,12X)//
175 FORMAT(1H1//////////1X,T40,'STRESS SIGMAXY-0 (PSI)')////
. 17X,6(I2,12X)//
210 FORMAT(1H,I12,(T15,6E14.5))
RETURN
END

```

```

SUBROUTINE TEMDIS (DELH,DELK,NX,NY,T0,T1,CURLIX,CURLIY,H)
DIMENSION T0(30,30),T1(30,30),CURLIX(30,30),CURLIY(30,30)

```

```

C
C THIS MODULE GIVES T0 AND T1 DISTRIBUTION FOR WHICH
C FOURTH ORDER POLYNOMIALS ARE ASSUMED.
C READ(5,10) U1,U2,U3,U4,U5,U6,U7,U8,U9,U10,U11,U12,U13,
. U14,U15
. READ(5,10) V1,V2,V3,V4,V5,V6,V7,V8,V9,V10,V11,V12,V13,
. V14,V15

```

```

DO 20 I=1,NX
DO 20 J=1,NY
U=I
V=J
X=(U-1.)*DELH
Y=(V-1.)*DELK
T0(I,J)=U1*X**4+U2*X**3+U3*X**2+U4*X*X**Y+U5*X*X**Y+U6*X*X
      +U7*X+U8*X**Y+U9*Y+U10*Y**Y+U11*X*Y**Y+U12*Y**3+U13*X*Y**3
      +U14*Y**4+CONST1
T1(I,J)=(V1*X**4+V2*X**3+V3*X**2+V4*X*X**Y+V5*X*X**Y+V6*X*X
      +V7*X+V8*X**Y+V9*Y+V10*Y**Y+V11*X*Y**Y+V12*Y**3+V13*X*Y**3
      +V14*Y**4+CONST2)*H
C
C   CURLIX, AND CURLY ARE FIRST DERIVATIVES OF T1 WITH
C   RESPECT TO X, AND Y, RESPECTIVELY.
CURLIX(I,J)=4.*V1*X**3+3.*V2*X*X**Y+3.*V3*X*X+2.*V4*X*Y
      +2.*V5*X*Y+2.*V6*X+V7+V8*Y+V11*Y**Y+V13*Y**3
CURLY(I,J)=V2*X**3+2.*V4*X*X**Y+V5*X*X+V9+2.*V10*Y+2.*V11
      *X*Y+3.*V12*Y**Y+3.*V13*X*Y**Y+4.*V14*Y**3
20 CONTINUE
WRITE(6,15) (J,J=1,6)
DO 30 I=1,NX
30 WRITE(6,60) I,(T0(I,J),J=1,NY)
   WRITE(6,23) (J,J=1,6)
DO 40 I=1,NX
40 WRITE(6,60) I,(T1(I,J),J=1,NY)
DO 70 I=1,NX
DO 70 J=1,NY
70 T1(I,J)=T1(I,J)/H
10 FORMAT(8F10,4)
15 FORMAT(1H1//////////1X,T35,TO TEMPERATURE DISTRIBUTION (DEGREES C
      )////////17X,6(12,11X)/)

```

```

23  FORMAT(1H1//////////1X,T15,'TEMPERATURE DISTRIBUTION DUE TO ZT1 ON
    . THE SURFACE (DEGREES C)'//17X,6(I2,11X)//)
60  FORMAT(1H,I12,(T15,6E13.4))
    RETURN
    END

```

```

SUBROUTINE DEFLXN (B,ANX,ANY,ANXY,NX,NY,DELH,DELK,W,H1,
    . ALFA1,ALFA2,ALFA3,ALFA6,ALFA6,TEMPW,TT1,CURLIX,CURLIY,ALFA)
    DIMENSION B(6,6),ANX(30,30),ANY(30,30),ANXY(30,30),ALFA(6)
    DIMENSION TT1(30,30),CURLIX(30,30),CURLIY(30,30)
    DIMENSION W(30,30),TEMPW(30,30),ALFA(6)
    X(Z) = (Z-1.)*DELH
    Y(Z) = (Z-1.)*DELK

```

C

```

    RIGHT HAND SIDE FUNCTION RHSW(X,Y)
    RHSW(X,Y) = ALFA1*(12.*V1*X*X+6.*(V2*X*Y+V3*X)+2.*(V4*Y*Y
    . +V5*Y+V6))+ALFA2*(2.*(V4*X*X+V10+V11*X)+6.*(V12*Y
    . +V13*X*Y)+12.*V14*Y*Y)+2.*ALFA6*(3.*V2*X*X+4.*V4*X*Y
    . +2.*V5*X+V8+2.*V11*Y+3.*V13*Y*Y)
    READ(5,10) V1,V2,V3,V4,V5,V6,V7,V8,V9,V10,V11,V12,V13,
    . V14
    EPSLNW=1.E-5
    NX4=NX-4
    NX3=NX-3
    NX2=NX-2
    NX1=NX-1
    NY4=NY-4
    NY3=NY-3
    NY2=NY-2
    NY1=NY-1

```

P1=B(1,1)/DELH\*\*4  
 Q1=-3.\*B(1,6)/(24.\*DELH\*\*3\*DELK)  
 R1=2.\*(B(1,2)+B(6,6))/(144.\*DELH\*\*2\*DELK\*\*2)  
 S1=-3.\*B(2,6)/(24.\*DELH\*DELK\*\*3)  
 T1=B(2,2)/DELK\*\*4

C THE COEFFICIENTS OF THE FINITE DIFFERENCE EQUATIONS FOR  
 C COMPUTING THE DEFLECTION.

DM1=Q1+R1+S1  
 DM2=-8.\*Q1-16.\*R1-2.\*S1  
 DM3=P1+30.\*R1  
 DM4=8.\*Q1-16.\*R1+2.\*S1  
 DM5=-Q1+R1-S1  
 DM6=-2.\*Q1-16.\*R1-8.\*S1  
 DM7=16.\*Q1+256.\*R1+16.\*S1  
 DM8=-4.\*P1-480.\*R1  
 DM9=-16.\*Q1+256.\*R1-16.\*S1  
 DM10=2.\*Q1-16.\*R1+8.\*S1  
 DM11=30.\*R1+T1  
 DM12=-480.\*R1-4.\*T1  
 DM13=6.\*P1+900.\*R1+6.\*T1  
 DM14=DM12  
 DM15=DM11  
 DM16=DM10  
 DM17=DM9  
 DM18=DM8  
 DM19=DM7  
 DM20=DM6  
 DM21=DM5  
 DM22=DM4  
 DM23=DM3  
 DM24=DM2



```

C
C
C
C
DM25=DM1
      COEFFICIENTS OF THE FINITE DIFFERENCE EQUATION
      REPRESENTING BOUNDARY CONDITION.
      B1= B(1,1)/DELH**3
      B2= -(B(1,6)/(DELH**2*DELK))
      B3= (B(1,2)-2.*B(6,6))/(DELH*DELK**2)
      B4= -2.*B(2,6)/DELK**3
      B11=2.*B(1,6)/DELH**3
      B22=(2.*B(1,6)-B(1,2))/(DELH**2*DELK)
      B33=B(2,6)/(DELH*DELK**2)
      B44=B(2,2)/DELK**3
      BT1=B(1,1)/DELH**2
      BT2=B(1,2)/DELK**2
      BT3=B(2,6)/(DELH*DELK)
      BT11=B(1,2)/DELH**2
      BT22=B(2,2)/DELK**2
      BT33=B(2,6)/(DELH*DELK)
C
C
      COMPUTE VALUES OF W ON BOUNDARIES.
      DO 20 I=1,NX
      DO 20 J=1,NY
      W(I,J)=0.
      TEMPW(I,J)=0.
      M=1
      DO 400 DO 100 J=1,NY3
      I=1
      TEMPW(I,J)= ((2.*BT2+BT3)*W(I,J+1)-BT2*W(I,J+2)+(2.*BT1+BT3)*
      . W(I+1,J)-BT3*W(I+1,J+1)-BT1*W(I+2,J)-ALFA(1)*TT1(I,J))/
      . (BT1+BT2+BT3)
      TEMPW(I+1,J)= ((B1+B2+B3+B4)*W(I,J)-(B2+2.*B3+3.*B4)*W(I,J+1)

```

```

    • +(B3+3.*B4)*W(I,J+2)-B4*W(I,J+3)+2.*(B2+B3)*W(I+1,J+1)
    • -B3*W(I+1,J+2)+(3.*B1+B2)*W(I+2,J)-B2*W(I+2,J+1)
    • -B1*W(I+3,J)-(ALFA(1)*CURL1X(I,J)-2.*ALFA(6)*
    • CURL1Y(I,J))/(3.*B1+2.*B2+B3)
    I=NX
    TEMPW(I,J)= ((2.*BT2-BT3)*W(I,J+1)-BT2*W(I,J+2)+(2.*BT1-BT3)*
    • W(I-1,J)+BT3*W(I-1,J+1)-BT1*W(I-2,J)-ALFA(1)*TT1(I,J))/
    • (BT1+BT2-BT3)
    TEMPW(I-1,J)= ((B1+B2+B3+B4)*W(I,J)-(B2+2.*B3+3.*B4)*W(I,J+1)
    • +(B3+3.*B4)*W(I,J+2)-B4*W(I,J+3)+2.*(B2+B3)*W(I-1,J+1)
    • -B3*W(I-1,J+2)+(3.*B1+B2)*W(I-2,J)-B2*W(I-2,J+1)
    • -B1*W(I-3,J)+(ALFA(1)*CURL1X(I,J)-2.*ALFA(6)*
    • CURL1Y(I,J))/(3.*B1+2.*B2+B3)
100 CONTINUE
    DO 120 J=NY2,NY
    I=1
    TEMPW(I,J)= ((2.*BT2-BT3)*W(I,J-1)-BT2*W(I,J-2)+(2.*BT1-BT3)*
    • W(I+1,J)+BT3*W(I+1,J-1)-BT1*W(I+2,J)-ALFA(1)*TT1(I,J))/
    • (BT1+BT2-BT3)
    TEMPW(I+1,J)= ((B1+B2+B3+B4)*W(I,J)-(B2+2.*B3+3.*B4)*W(I,J-1)
    • +(B3+3.*B4)*W(I,J-2)-B4*W(I,J-3)+2.*(B2+B3)*W(I+1,J-1)
    • -B3*W(I+1,J-2)+(3.*B1+B2)*W(I+2,J)-B2*W(I+2,J-1)
    • -B1*W(I+3,J)+(ALFA(1)*CURL1X(I,J)-2.*ALFA(6)*
    • CURL1Y(I,J))/(3.*B1+2.*B2+B3)
    I=NX
    TEMPW(I,J)= ((2.*BT2+BT3)*W(I,J-1)-BT2*W(I,J-2)+(2.*BT1+BT3)*
    • W(I-1,J)-BT3*W(I-1,J-1)-BT1*W(I-2,J)-ALFA(1)*TT1(I,J))/
    • (BT1+BT2+BT3)
    TEMPW(I-1,J)= ((B1+B2+B3+B4)*W(I,J)-(B2+2.*B3+3.*B4)*W(I,J-1)
    • +(B3+3.*B4)*W(I,J-2)-B4*W(I,J-3)+2.*(B2+B3)*W(I-1,J-1)
    • -B3*W(I-1,J-2)+(3.*B1+B2)*W(I-2,J)-B2*W(I-2,J-1)
    • -B1*W(I-3,J)-(ALFA(1)*CURL1X(I,J)-2.*ALFA(6)*

```

```

120      CURL1Y(I,J))/(3.*B1+2.*B2+B3)
      CONTINUE
      DO 150 I=2,NX3
      J=1
      TEMPW(I,J)= ((2.*BT22+BT33)*W(I,J+1)-BT22*W(I,J+2)+(2.*BT11+
      BT33)*W(I+1,J)-BT33*W(I+1,J+1)-BT11*W(I+2,J)-ALFA(2)*
      TT1(I,J))/(BT11+BT22+BT33)
      TEMPW(I,J+1)= ((B11+B22+B33+B44)*W(I,J)+(B33+3.*B44)*W(I,J+2)
      -B44*W(I,J+3)-(3.*B11+2.*B22+B33)*W(I+1,J)+
      2.*(B22+B33)*W(I+1,J+1)-B33*W(I+1,J+2)+(3.*B11+B22)*
      W(I+2,J)-B22*W(I+2,J+1)-B11*W(I+3,J)-(2.*ALFA(6)*
      CURL1X(I,J)-ALFA(2)*CURL1Y(I,J))/(B22+2.*B33+3.*B44)
      J=NY
      TEMPW(I,J)= ((2.*BT22-BT33)*W(I,J-1)-BT22*W(I,J-2)+(2.*BT11-
      BT33)*W(I+1,J)+BT33*W(I+1,J-1)-BT11*W(I+2,J)-ALFA(2)*
      TT1(I,J))/(BT11+BT22-BT33)
      TEMPW(I,J-1)= ((B11+B22+B33+B44)*W(I,J)+(B33+3.*B44)*W(I,J-2)
      -B44*W(I,J-3)-(3.*B11+2.*B22+B33)*W(I+1,J)+
      2.*(B22+B33)*W(I+1,J-1)-B33*W(I+1,J-2)+(3.*B11+B22)*
      W(I+2,J)-B22*W(I+2,J-1)-B11*W(I+3,J)+(2.*ALFA(6)*
      CURL1X(I,J)-ALFA(2)*CURL1Y(I,J))/(B22+2.*B33+3.*B44)
      CONTINUE
      DO 160 I=NX2,NX1
      J=1
      TEMPW(I,J)= ((2.*BT22-BT33)*W(I,J+1)-BT22*W(I,J+2)+(2.*BT11-
      BT33)*W(I-1,J)+BT33*W(I-1,J+1)-BT11*W(I-2,J)-ALFA(2)*
      TT1(I,J))/(BT11+BT22-BT33)
      TEMPW(I,J+1)= ((B11+B22+B33+B44)*W(I,J)+(B33+3.*B44)*W(I,J+2)
      -B44*W(I,J+3)-(3.*B11+2.*B22+B33)*W(I-1,J)+
      2.*(B22+B33)*W(I-1,J+1)-B33*W(I-1,J+2)+(3.*B11+B22)*
      W(I-2,J)-B22*W(I-2,J+1)-B11*W(I-3,J)+(2.*ALFA(6)*
      CURL1X(I,J)-ALFA(2)*CURL1Y(I,J))/(B22+2.*B33+3.*B44)

```

```

J=NY
TEMPW(I,J) = ((2.*BT22+BT33)*W(I,J-1)-BT22*W(I,J-2)+(2.*BT11+
  BT33)*W(I-1,J)-BT33*W(I-1,J-1)-BT11*W(I-2,J)-ALFA(2)*
  TT1(I,J))/(BT11+BT22+BT33)
TEMPW(I,J-1) = ((B11+B22+B33+B44)*W(I,J)+(B33+3.*B44)*W(I,J-2)
  -B44*W(I,J-3)-(3.*B11+2.*B22+B33)*W(I-1,J)+
  2.*(B22+B33)*W(I-1,J-1)-B33*W(I-1,J-2)+(3.*B11+B22)*
  W(I-2,J)-B22*W(I-2,J-1)-B11*W(I-3,J)-(2.*ALFA(6)*
  CURLIX(I,J)-ALFA(2)*CURLIY(I,J)))/(B22+2.*B33+3.*B44)
160 CONTINUE

```

C  
C COMPUTE DEFLECTION AT THE INNER POINTS.  
C

```

DO 350 J=3,NY2
DO 350 I=3,NX2
U=I
V=J

```

```

350 TEMPW(I,J) = -(DM1*W(I-2,J-2)+DM2*W(I-2,J-1)+DM3*W(I-2,J)
  +DM4*W(I-2,J+1)+DM5*W(I-2,J+2)+DM6*W(I-1,J-2)+DM7*
  W(I-1,J-1)+(DM8-H1*ANX(I,J))/(DELH**2))*W(I-1,J)+DM9*
  W(I-1,J+1)+DM10*W(I-1,J+2)+DM11*W(I,J-2)+(DM12-H1*
  ANY(I,J))/(DELK**2))*W(I,J-1)+(DM14-H1*ANY(I,J)/
  (DELK**2))*W(I,J+1)+DM15*W(I,J+2)+DM16*W(I+1,J-2)+DM17
  *W(I+1,J-1)+(DM18-H1*ANX(I,J))/(DELH**2))*W(I+1,J)+
  DM19*W(I+1,J+1)+DM20*W(I+1,J+2)+DM21*W(I+2,J-2)+DM22*
  W(I+2,J-1)+DM23*W(I+2,J)+DM24*W(I+2,J+1)+DM25*W(I+2,J+2)
  -H1*ANXY(I,J)*W(I-1,J-1)-W(I-1,J+1)-W(I+1,J-1)+
  W(I+1,J+1))/(2.*DELH*DELK)+RHSW(X(U),Y(V))/(DM13+H1*
  2.*(ANX(I,J)/(DELH**2))+ANY(I,J)/(DELK**2)))

```

C  
C CHECK FOR CONVERGENCE OF W.  
C DO 370 I=1,NX

```

DO 370 J=1,NY
IF(ABS(TEMPW(I,J)-W(I,J)).GT.EPSLNW) GO TO 375
370 CONTINUE
375 M=M+1
C
C IF NO CONVERGENCE REPLACE OLD VALUES OF W WITH NEW
C VALUES OF W, THAT IS TEMPW.
DO 390 I=1,NX
DO 390 J=1,NY
W(I,J)=TEMPW(I,J)
IF (M.GT.40) GO TO 380
GO TO 400
WRITE (6,440) (J,J=1,6)
DO 395 I=1,NX
395 WRITE (6,500) I,(W(I,J),J=1,NY)
10 FORMAT(8F10.4)
440 FORMAT (1H1//////////1H,T35,'DEFLECTION (INCH)')////
17X,6(I2,11X)
500 FORMAT (1H,I12,(T15,6E13.4))
RETURN
END
C
C SUBROUTINE SIGMA1(W,B,DELH,DELK,H,NX,NY,TELX,TELY,TELXY,
C ALFA,T1)
C DIMENSION W(30,30),TELX(30,30),TELY(30,30),TELXY(30,30)
C DIMENSION T1(30,30),B(6,6),ALFA(6)
C TELX, TELY, AND TELXY COSTITUTE THE STRESS TENSOR TAU-1.
NX4=NX-4
NX3=NX-3

```

```

NX2=NX-2
NX1=NX-1
NY4=NY-4
NY3=NY-3
NY2=NY-2
NY1=NY-1
C
C
COEFFICIENTS OF THE FINITE DIFFERENCE EQUATIONS FOR
STRESSES TAU-1.
P2=B(1,1)/(12.*DELH**2.)
Q2=B(1,2)/(12.*DELK**2.)
R2=B(1,6)/(4.*DELH*DELK)
P3=B(1,2)/(12.*DELH**2.)
Q3=B(2,2)/(12.*DELK**2.)
R3=B(2,6)/(4.*DELH*DELK)
P4=B(1,6)/(12.*DELH**2.)
Q4=B(2,6)/(12.*DELK**2.)
R4=B(6,6)/(4.*DELH*DELK)
DO 800 I=1,NX4

```

```

C
C
STRESS COMPUTATION.
DO 800 J=1,4
TELX(I,J)=-H*((35.*(P2+Q2)+9.*R2)*W(I,J)-(104.*Q2+12.*R2)
* W(I,J+1)+(114.*Q2+3.*R2)*W(I,J+2)-56.*Q2*W(I,J+3)+
* 11.*Q2*W(I,J+4)-(104.*P2+12.*R2)*W(I+1,J)+16.*R2*W(I+1,J+1)
* -4.*R2*W(I+1,J+2)+(114.*P2+3.*R2)*W(I+2,J)-4.*R2*
* W(I+2,J+1)+R2*W(I+2,J+2)-56.*P2*W(I+3,J)+11.*P2*W(I+4,J))
TELY(I,J)=-H*((35.*(P3+Q3)+9.*R3)*W(I,J+2)-56.*Q3*W(I,J+3)+
* W(I,J+1)+(114.*Q3+3.*R3)*W(I,J+2)-56.*R3*W(I+1,J+1)
* 11.*Q3*W(I,J+4)-(104.*P3+12.*R3)*W(I+1,J)+16.*R3*
* -4.*R3*W(I+1,J+2)+(114.*P3+3.*R3)*W(I+2,J)-4.*R3*
* W(I+2,J+1)+R3*W(I+2,J+2)-56.*P3*W(I+3,J)+11.*P3*W(I+4,J))
TELY(I,J)=-H*((35.*(P4+Q4)+9.*R4)*W(I,J)-(104.*Q4+12.*R4)

```

```

    • #W(I,J+1)+(114.*Q4+3.*R4)*W(I,J+2)-56.*Q4*W(I,J+3)+
    • 11.*Q4*W(I,J+4)-(104.*P4+12.*R4)*W(I+1,J)+16.*R4*W(I+1,J+1)
    • -4.*R4*W(I+1,J+2)+(114.*P4+3.*R4)*W(I+2,J)-4.*R4*
    • W(I+2,J+1)+R4*W(I+2,J+2)-56.*P4*W(I+3,J)+11.*P4*W(I+4,J))

```

800 CONTINUE

DO 810 I=1,4

DO 810 J=5,NY

```

    TELX(I,J)=-H*((35.*(P2+Q2)-9.*R2)*W(I,J)-(104.*Q2-12.*R2)
    • #W(I,J-1)+(114.*Q2-3.*R2)*W(I,J-2)-56.*Q2*W(I,J-3)+
    • 11.*Q2*W(I,J-4)-(104.*P2-12.*R2)*W(I+1,J)-16.*R2*W(I+1,J-1)
    • +4.*R2*W(I+1,J-2)+(114.*P2-3.*R2)*W(I+2,J)+4.*R2*
    • W(I+2,J-1)-R2*W(I+2,J-2)-56.*P2*W(I+3,J)+11.*P2*W(I+4,J))
    TELY(I,J)=-H*((35.*(P3+Q3)-9.*R3)*W(I,J)-(104.*Q3-12.*R3)
    • #W(I,J-1)+(114.*Q3-3.*R3)*W(I,J-2)-56.*Q3*W(I,J-3)+
    • 11.*Q3*W(I,J-4)-(104.*P3-12.*R3)*W(I+1,J)-16.*R3*W(I+1,J-1)
    • +4.*R3*W(I+1,J-2)+(114.*P3-3.*R3)*W(I+2,J)+4.*R3*
    • W(I+2,J-1)-R3*W(I+2,J-2)-56.*P3*W(I+3,J)+11.*P3*W(I+4,J))
    TELXY(I,J)=-H*((35.*(P4+Q4)-9.*R4)*W(I,J)-(104.*Q4-12.*R4)
    • #W(I,J-1)+(114.*Q4-3.*R4)*W(I,J-2)-56.*Q4*W(I,J-3)+
    • 11.*Q4*W(I,J-4)-(104.*P4-12.*R4)*W(I+1,J)-16.*R4*W(I+1,J-1)
    • +4.*R4*W(I+1,J-2)+(114.*P4-3.*R4)*W(I+2,J)+4.*R4*
    • W(I+2,J-1)-R4*W(I+2,J-2)-56.*P4*W(I+3,J)+11.*P4*W(I+4,J))

```

810 CONTINUE

DO 820 I=5,NX

DO 820 J=NY3,NY

```

    TELX(I,J)=-H*((35.*(P2+Q2)+9.*R2)*W(I,J)-(104.*Q2+12.*R2)
    • #W(I,J-1)+(114.*Q2+3.*R2)*W(I,J-2)-56.*Q2*W(I,J-3)+
    • 11.*Q2*W(I,J-4)-(104.*P2+12.*R2)*W(I-1,J)+16.*R2*W(I-1,J-1)
    • -4.*R2*W(I-1,J-2)+(114.*P2+3.*R2)*W(I-2,J)-4.*R2*
    • W(I-2,J-1)+R2*W(I-2,J-2)-56.*P2*W(I-3,J)+11.*P2*W(I-4,J))
    TELY(I,J)=-H*((35.*(P3+Q3)+9.*R3)*W(I,J)-(104.*Q3+12.*R3)
    • #W(I,J-1)+(114.*Q3+3.*R3)*W(I,J-2)-56.*Q3*W(I,J-3)+

```

```

    • 11.*Q3*W(I,J-4)-(104.*P3+12.*R3)*W(I-1,J)+16.*R3*W(I-1,J-1)
    • -4.*R3*W(I-1,J-2)+(114.*P3+3.*R3)*W(I-2,J)-4.*R3*
    • W(I-2,J-1)+R3*W(I-2,J-2)-56.*P3*W(I-3,J)+11.*P3*W(I-4,J))
    TELXY(I,J)=-H*((35.*(P4+Q4)+9.*R4)*W(I,J)-(104.*Q4+12.*R4)
    • *W(I,J-1)+(114.*Q4+3.*R4)*W(I,J-2)-56.*Q4*W(I,J-3)+
    • 11.*Q4*W(I,J-4)-(104.*P4+12.*R4)*W(I-1,J)+16.*R4*W(I-1,J-1)
    • -4.*R4*W(I-1,J-2)+(114.*P4+3.*R4)*W(I-2,J)-4.*R4*
    • W(I-2,J-1)+R4*W(I-2,J-2)-56.*P4*W(I-3,J)+11.*P4*W(I-4,J))

```

820 CONTINUE

DO 830 I=NX3,NX

DO 830 J=1,NY4

```

    TELX(I,J)=-H*((35.*(P2+Q2)-9.*R2)*W(I,J)-(104.*Q2-12.*R2)
    • *W(I,J+1)+(114.*Q2-3.*R2)*W(I,J+2)-56.*Q2*W(I,J+3)+
    • 11.*Q2*W(I,J+4)-(104.*P2-12.*R2)*W(I-1,J)-16.*R2*W(I-1,J+1)
    • +4.*R2*W(I-1,J+2)+(114.*P2-3.*R2)*W(I-2,J)+4.*R2*
    • W(I-2,J+1)-R2*W(I-2,J+2)-56.*P2*W(I-3,J)+11.*P2*W(I-4,J))
    TELY(I,J)=-H*((35.*(P3+Q3)-9.*R3)*W(I,J)-(104.*Q3-12.*R3)
    • *W(I,J+1)+(114.*Q3-3.*R3)*W(I,J+2)-56.*Q3*W(I,J+3)+
    • 11.*Q3*W(I,J+4)-(104.*P3-12.*R3)*W(I-1,J)-16.*R3*W(I-1,J+1)
    • +4.*R3*W(I-1,J+2)+(114.*P3-3.*R3)*W(I-2,J)+4.*R3*
    • W(I-2,J+1)-R3*W(I-2,J+2)-56.*P3*W(I-3,J)+11.*P3*W(I-4,J))
    TELXY(I,J)=-H*((35.*(P4+Q4)-9.*R4)*W(I,J)-(104.*Q4-12.*R4)
    • *W(I,J+1)+(114.*Q4-3.*R4)*W(I,J+2)-56.*Q4*W(I,J+3)+
    • 11.*Q4*W(I,J+4)-(104.*P4-12.*R4)*W(I-1,J)-16.*R4*W(I-1,J+1)
    • +4.*R4*W(I-1,J+2)+(114.*P4-3.*R4)*W(I-2,J)+4.*R4*
    • W(I-2,J+1)-R4*W(I-2,J+2)-56.*P4*W(I-3,J)+11.*P4*W(I-4,J))

```

830 CONTINUE

DO 840 I=5,NX4

DO 840 J=5,NY4

```

    TELX(I,J)=-H*(-P2*W(I-2,J)+16.*P2*W(I-1,J)+R2*(W(I-1,J-1)-
    • W(I-1,J+1))-30.*(P2+Q2)*W(I,J)-Q2*W(I,J-2)+16.*Q2*W(I,J-1)-
    • R2*(W(I+1,J-1)-W(I+1,J+1))+16.*(P2*W(I+1,J)+Q2*W(I,J+1))

```



```

      • -P2*W(I+2,J)-Q2*W(I,J+2)
      TELY(I,J)=-H*(-P3*W(I-2,J)+16.*P3*W(I-1,J)+R3*(W(I-1,J-1))-
      • W(I-1,J+1))-30.*(P3+Q3)*W(I,J)-Q3*W(I,J-2)+16.*Q3*W(I,J-1)-
      • R3*(W(I+1,J-1))-W(I+1,J+1))+16.*(P3*W(I+1,J)+Q3*W(I,J+1))
      • -P3*W(I+2,J)-Q3*W(I,J+2)
      TELY(I,J)=-H*(-P4*W(I-2,J)+16.*P4*W(I-1,J)+R4*(W(I-1,J-1))-
      • W(I-1,J+1))-30.*(P4+Q4)*W(I,J)-Q4*W(I,J-2)+16.*Q4*W(I,J-1)-
      • R4*(W(I+1,J-1))-W(I+1,J+1))+16.*(P4*W(I+1,J)+Q4*W(I,J+1))
      • -P4*W(I+2,J)-Q4*W(I,J+2)
840  CONTINUE
      ALFA1H=H*ALFA(1)
      ALFA2H=H*ALFA(2)
      ALFA6H=H*ALFA(6)
      DO 50 I=1,NX
      DO 50 J=1,NY
      TELY(I,J)=TELY(I,J)-ALFA1H*T1(I,J)
      TELY(I,J)=TELY(I,J)-ALFA2H*T1(I,J)
      TELY(I,J)=TELY(I,J)-ALFA6H*T1(I,J)
50
C
C
C      PRINT STRESS DISTRIBUTION TAU-1 ON THE UPPER SURFACE
      WRITE (6,55) (J,J=1,6)
      DO 60 I=1,NX
      WRITE (6,1100) I,(TELY(I,J),J=1,NY)
      WRITE (6,65) (J,J=1,6)
      DO 70 I=1,NX
      WRITE (6,1100) I,(TELY(I,J),J=1,NY)
      WRITE (6,75) (J,J=1,6)
      DO 80 I=1,NX
      WRITE (6,1100) I,(TELY(I,J),J=1,NY)
80
55  FORMAT(1H1//////////1X,T35,'STRESS SIGMAX-1 ON THE UPPER SURFACE (
      •PSI)')//17X,6(I2,12X)//

```

```

65 FORMAT(1H1//////////1X,T35,STRESS SIGMAY-1 ON THE UPPER SURFACE (
    .PSI)////17X,6(12,12X)//
75 FORMAT(1H1//////////1X,T35,STRESS SIGMAXY-1 ON THE UPPER SURFACE
    .(PSI)////17X,6(12,12X)//
1100 FORMAT (1H,112,(T15,6E14.5))
RETURN
END

```

INPUT DATA

C NX,NY,DELH,DELK,H

15 11 .4 .4 .1

C ELASTIC CONSTANTS (COMPLIANCES) A(I,J)\*1.E-8

2.	-5.	-4.5	0.	0.	0.	-7.
-5.	5.	-2.2	0.	0.	0.	10.
-4.5	-2.2	10.2	0.	0.	0.	9.
0.	0.	0.	12.	13.75	0.	0.
0.	0.	0.	13.75	11.	0.	0.
-7.	10.	9.	0.	0.	0.	10.

C LINEAR COEFFICIENTS OF THERMAL EXPANSION (ALFA\*1.E-6)

12.5 10. 9. 9.5

C POLYNOMIAL COEFFICIENTS U1,U2,....U15, AND V1,V2,....V15



## T0 TEMPERATURE DISTRIBUTION (DEGREES C)

	1	2	3	4	5	6
1	0.1000E+03	0.1000E+03	0.1003E+03	0.1010E+03	0.1025E+03	0.1048E+03
2	0.1083E+03	0.1132E+03	0.1197E+03	0.1280E+03	0.1384E+03	0.1500E+03
3	0.1003E+03	0.1005E+03	0.1009E+03	0.1018E+03	0.1034E+03	0.1059E+03
4	0.1096E+03	0.1146E+03	0.1213E+03	0.1298E+03	0.1403E+03	0.1500E+03
5	0.1013E+03	0.1016E+03	0.1022E+03	0.1033E+03	0.1050E+03	0.1077E+03
6	0.1115E+03	0.1167E+03	0.1235E+03	0.1322E+03	0.1429E+03	0.1500E+03
7	0.1029E+03	0.1034E+03	0.1041E+03	0.1054E+03	0.1073E+03	0.1101E+03
8	0.1141E+03	0.1194E+03	0.1264E+03	0.1352E+03	0.1461E+03	0.1500E+03
9	0.1051E+03	0.1058E+03	0.1067E+03	0.1081E+03	0.1101E+03	0.1131E+03
10	0.1173E+03	0.1228E+03	0.1299E+03	0.1389E+03	0.1499E+03	0.1500E+03
11	0.1080E+03	0.1088E+03	0.1099E+03	0.1114E+03	0.1137E+03	0.1168E+03
12	0.1211E+03	0.1268E+03	0.1341E+03	0.1432E+03	0.1544E+03	0.1500E+03
13	0.1115E+03	0.1125E+03	0.1137E+03	0.1154E+03	0.1178E+03	0.1211E+03
14	0.1256E+03	0.1314E+03	0.1389E+03	0.1482E+03	0.1595E+03	0.1500E+03
15	0.1157E+03	0.1168E+03	0.1182E+03	0.1201E+03	0.1226E+03	0.1261E+03
16	0.1307E+03	0.1367E+03	0.1443E+03	0.1538E+03	0.1653E+03	0.1500E+03
17	0.1205E+03	0.1218E+03	0.1233E+03	0.1254E+03	0.1281E+03	0.1317E+03
18	0.1365E+03	0.1426E+03	0.1504E+03	0.1600E+03	0.1717E+03	0.1500E+03
19	0.1259E+03	0.1274E+03	0.1291E+03	0.1313E+03	0.1341E+03	0.1379E+03
20	0.1429E+03	0.1492E+03	0.1571E+03	0.1669E+03	0.1787E+03	0.1500E+03
21	0.1320E+03	0.1336E+03	0.1355E+03	0.1378E+03	0.1409E+03	0.1448E+03
22	0.1499E+03	0.1564E+03	0.1645E+03	0.1744E+03	0.1864E+03	0.1500E+03
23	0.1387E+03	0.1405E+03	0.1425E+03	0.1450E+03	0.1482E+03	0.1523E+03
24	0.1576E+03	0.1642E+03	0.1725E+03	0.1826E+03	0.1947E+03	0.1500E+03
25	0.1461E+03	0.1480E+03	0.1502E+03	0.1529E+03	0.1562E+03	0.1605E+03
26	0.1659E+03	0.1727E+03	0.1811E+03	0.1914E+03	0.2037E+03	0.1500E+03
27	0.1541E+03	0.1562E+03	0.1585E+03	0.1614E+03	0.1649E+03	0.1693E+03
28	0.1749E+03	0.1818E+03	0.1904E+03	0.2008E+03	0.2133E+03	0.1500E+03
29	0.1627E+03	0.1650E+03	0.1675E+03	0.1705E+03	0.1741E+03	0.1787E+03
30	0.1845E+03	0.1916E+03	0.2003E+03	0.2109E+03	0.2235E+03	0.1500E+03

## APPENDIX B

## RESULTS

TEMPERATURE DISTRIBUTION DUE TO ZT1 ON THE SURFACE (DEGREES C)

	1	2	3	4	5	6
1	0.4000E+02	0.4000E+02	0.4000E+02	0.4000E+02	0.4000E+02	0.4000E+02
2	0.4000E+02	0.4000E+02	0.4011E+02	0.4019E+02	0.4028E+02	0.4041E+02
3	0.4058E+02	0.4079E+02	0.4106E+02	0.4139E+02	0.4179E+02	0.4090E+02
4	0.4000E+02	0.4012E+02	0.4025E+02	0.4042E+02	0.4063E+02	0.4090E+02
5	0.4125E+02	0.4169E+02	0.4224E+02	0.4291E+02	0.4372E+02	0.4155E+02
6	0.4000E+02	0.4022E+02	0.4046E+02	0.4074E+02	0.4110E+02	0.4155E+02
7	0.4211E+02	0.4281E+02	0.4367E+02	0.4472E+02	0.4597E+02	0.4242E+02
8	0.4000E+02	0.4036E+02	0.4075E+02	0.4121E+02	0.4175E+02	0.4242E+02
9	0.4324E+02	0.4425E+02	0.4547E+02	0.4694E+02	0.4868E+02	0.4360E+02
10	0.4000E+02	0.4057E+02	0.4117E+02	0.4185E+02	0.4265E+02	0.4360E+02
11	0.4474E+02	0.4612E+02	0.4776E+02	0.4971E+02	0.5200E+02	0.4516E+02
12	0.4000E+02	0.4085E+02	0.4174E+02	0.4273E+02	0.4386E+02	0.4516E+02
13	0.4670E+02	0.4852E+02	0.5066E+02	0.5317E+02	0.5609E+02	0.4719E+02
14	0.4000E+02	0.4122E+02	0.4250E+02	0.4388E+02	0.4543E+02	0.4719E+02
15	0.4922E+02	0.5157E+02	0.5430E+02	0.5746E+02	0.6110E+02	0.4975E+02
16	0.4000E+02	0.4170E+02	0.4347E+02	0.4536E+02	0.4743E+02	0.4975E+02
17	0.5238E+02	0.5538E+02	0.5880E+02	0.6272E+02	0.6719E+02	0.5293E+02
18	0.4000E+02	0.4231E+02	0.4469E+02	0.4721E+02	0.4993E+02	0.5293E+02
19	0.5628E+02	0.6004E+02	0.6428E+02	0.6908E+02	0.7450E+02	0.5680E+02
20	0.4000E+02	0.4305E+02	0.4618E+02	0.4947E+02	0.5298E+02	0.5680E+02
21	0.6100E+02	0.6567E+02	0.7087E+02	0.7669E+02	0.8320E+02	0.6144E+02
22	0.4000E+02	0.4395E+02	0.4798E+02	0.5219E+02	0.5664E+02	0.6144E+02
23	0.6665E+02	0.7238E+02	0.7869E+02	0.8568E+02	0.9343E+02	0.6692E+02
24	0.4000E+02	0.4502E+02	0.5012E+02	0.5541E+02	0.6098E+02	0.6692E+02
25	0.7332E+02	0.8027E+02	0.8786E+02	0.9619E+02	0.1054E+03	0.7332E+02
26	0.4000E+02	0.4626E+02	0.5263E+02	0.5919E+02	0.6606E+02	0.7332E+02
27	0.8108E+02	0.8945E+02	0.9851E+02	0.1084E+03	0.1191E+03	0.8072E+02
28	0.4000E+02	0.4771E+02	0.5554E+02	0.6357E+02	0.7193E+02	0.8072E+02
29	0.9005E+02	0.1000E+03	0.1107E+03	0.1223E+03	0.1349E+03	0.8072E+02

STRESS SIGMAX-0 (PSI)

	1	2	3	4	5	6
1	0.0	0.0	0.0	0.0	0.0	0.0
2	0.0	0.0	0.0	0.0	0.0	0.0
3	0.0	0.0	0.0	0.0	0.0	0.0
4	0.99754E+00	-0.25646E+03	0.44638E+03	-0.64536E+03	0.44638E+03	-0.64536E+03
5	0.57660E+03	-0.26318E+03	0.34983E+03	0.11925E+03	-0.82806E+03	0.82806E+03
6	-0.37022E+03	0.58296E+03	-0.83500E+03	0.89916E+03	-0.83500E+03	0.89916E+03
7	-0.97755E+03	0.10938E+04	-0.56630E+02	0.13727E+03	-0.47381E+03	0.47381E+03
8	-0.14631E+03	-0.53939E+03	0.10679E+04	-0.13312E+04	-0.435468E+03	0.435468E+03
9	-0.28707E+03	-0.10401E+04	0.10867E+04	-0.16457E+03	-0.50961E+03	0.50961E+03
10	-0.54690E+03	0.72972E+03	-0.98931E+03	0.14781E+04	0.39141E+03	-0.47894E+03
11	0.51682E+03	0.20272E+04	-0.82560E+03	0.59091E+03	-0.66192E+03	0.66192E+03
12	-0.18633E+03	-0.55495E+03	0.12974E+04	-0.15716E+04	-0.41974E+03	0.41974E+03
13	-0.50787E+03	-0.18937E+04	0.20754E+04	-0.55985E+03	-0.49110E+03	0.49110E+03
14	-0.40661E+03	0.52054E+03	-0.91910E+03	0.15003E+04	0.39854E+03	-0.52922E+03
15	0.61937E+03	0.24093E+04	-0.14426E+04	0.12344E+04	-0.90754E+03	0.90754E+03
16	-0.16284E+03	-0.45177E+03	0.99114E+03	-0.12497E+04	-0.33560E+03	0.33560E+03
17	-0.54355E+03	-0.20285E+04	0.25651E+04	-0.87680E+03	-0.43635E+03	0.43635E+03
18	-0.28313E+03	0.21947E+03	-0.59347E+03	0.11357E+04	0.29915E+03	-0.38948E+03
19	0.48339E+03	0.18643E+04	-0.14477E+04	0.17330E+04	-0.11733E+04	0.11733E+04
20	-0.18742E+03	-0.21826E+03	0.47649E+03	-0.67632E+03	-0.18381E+03	0.18381E+03
21	-0.36979E+03	-0.13654E+04	0.19876E+04	-0.66395E+03	-0.43643E+03	0.43643E+03
22	-0.20213E+03	0.98263E+01	-0.18920E+03	0.57175E+03	-0.58898E+03	0.58898E+03
23	-0.96427E+03	0.90553E+03	-0.96427E+03	0.13926E+04	-0.93512E+03	0.93512E+03
24	-0.15147E+03	-0.66679E+01	0.14900E+03	-0.16228E+03	0.39175E+03	-0.46261E+03
25	0.63856E+03	-0.46261E+03	0.63856E+03	-0.28049E+03	-0.11580E+03	0.11580E+03
26	0.0	0.0	0.0	0.0	0.0	0.0
27	0.0	0.0	0.0	0.0	0.0	0.0
28	0.0	0.0	0.0	0.0	0.0	0.0
29	0.0	0.0	0.0	0.0	0.0	0.0

STRESS SIGMAY-0 (PSI)

	1	2	3	4	5	6
1	0.0	0.0	0.99754E+00	-0.25447E+03	-0.63553E+02	-0.51800E+03
2	-0.39584E+03	-0.53687E+03	-0.32806E+03	0.0	0.0	0.53314E+03
3	0.0	0.0	-0.37222E+03	0.35145E+03	-0.65264E+03	0.0
4	-0.41183E+03	0.26339E+03	0.18231E+03	0.0	0.0	-0.91108E+03
5	0.0	0.0	0.59513E+03	-0.77151E+03	0.10461E+04	0.0
6	0.12399E+04	-0.99968E+02	0.10995E+03	0.0	0.0	0.17958E+04
7	0.0	0.0	-0.62449E+03	0.11425E+04	-0.10507E+04	0.0
8	-0.12906E+04	0.82429E+03	-0.11652E+03	0.0	0.0	0.46399E+03
9	0.0	0.0	0.76115E+03	-0.10315E+04	-0.28792E+03	0.0
10	-0.35463E+03	-0.99968E+02	0.10995E+03	0.0	0.0	0.47828E+03
11	0.0	0.0	-0.58085E+03	0.11985E+04	0.39756E+03	0.0
12	0.51941E+03	0.82429E+03	-0.11652E+03	0.0	0.0	0.49756E+03
13	0.0	0.0	0.46405E+03	-0.11197E+04	-0.40500E+03	0.0
14	-0.60695E+03	-0.12600E+04	0.32314E+03	0.0	0.0	0.49118E+03
15	0.0	0.0	-0.36406E+03	0.91542E+03	0.37919E+03	0.0
16	0.62965E+03	0.17705E+04	-0.58726E+03	0.0	0.0	0.43147E+03
17	0.0	0.0	0.21601E+03	-0.67696E+03	-0.34512E+03	0.0
18	-0.58457E+03	-0.21302E+04	0.88763E+03	0.0	0.0	0.36136E+03
19	0.0	0.0	-0.11043E+03	0.44497E+03	0.28701E+03	0.0
20	0.48960E+03	0.23047E+04	-0.12081E+04	0.0	0.0	0.29703E+03
21	0.0	0.0	0.65381E+02	-0.11382E+03	-0.20149E+03	0.0
22	-0.32537E+03	-0.20591E+04	0.14738E+04	0.0	0.0	0.11446E+04
23	0.0	0.0	-0.11043E+03	0.44497E+03	-0.73529E+03	0.0
24	-0.11869E+04	0.19823E+04	-0.12356E+04	0.0	0.0	0.4654E+03
25	0.0	0.0	0.65381E+02	-0.11382E+03	0.71087E+03	0.0
26	0.10495E+04	-0.10936E+04	0.13180E+04	0.0	0.0	0.39460E+03
27	0.0	0.0	0.10080E+03	0.22476E+03	-0.13848E+03	0.0
28	-0.44481E+03	0.54653E+03	-0.70352E+03	0.0	0.0	0.49015E+03
29	0.0	0.0	-0.15147E+03	-0.30960E+03	-0.31874E+03	0.0
30	-0.26982E+03	-0.51209E+03	-0.11580E+03	0.0	0.0	0.0

STRESS SIGMAXY-0 (PSI)

	1	2	3	4	5	6
1	0.0	0.0	0.0	0.0	0.0	0.0
2	0.0	0.0	0.0	0.0	0.0	0.0
3	0.0	0.0	0.0	0.0	0.0	0.0
4	0.0	0.0	0.0	0.0	0.0	0.0
5	0.0	0.0	0.0	0.0	0.0	0.0
6	0.0	0.0	0.0	0.0	0.0	0.0
7	0.0	0.0	0.0	0.0	0.0	0.0
8	0.0	0.0	0.0	0.0	0.0	0.0
9	0.0	0.0	0.0	0.0	0.0	0.0
10	0.0	0.0	0.0	0.0	0.0	0.0
11	0.0	0.0	0.0	0.0	0.0	0.0
12	0.0	0.0	0.0	0.0	0.0	0.0
13	0.0	0.0	0.0	0.0	0.0	0.0
14	0.0	0.0	0.0	0.0	0.0	0.0
15	0.0	0.0	0.0	0.0	0.0	0.0



DEFLECTION (INCH)

	1	2	3	4	5	6
1	-0.3923E-02	-0.1137E-01	-0.1878E-01	-0.1886E-01	-0.1610E-01	-0.3522E-01
2	-0.8289E-02	-0.9520E-02	0.1319E-02	-0.7740E-03	-0.4542E-03	
3	0.9924E-02	0.7324E-03	0.1875E-02	0.1189E-02	0.3851E-02	0.1407E-02
4	0.3028E-02	0.3366E-04	0.1054E-02	-0.4684E-04	0.2669E-03	
5	0.1631E-01	0.1820E-02	-0.1082E-02	-0.7828E-03	-0.9031E-03	-0.1129E-02
6	-0.8721E-03	-0.7408E-04	-0.2369E-04	-0.1476E-03	-0.3404E-03	
7	0.1472E-01	0.1517E-02	0.1858E-03	0.2287E-03	0.3036E-03	0.3287E-03
8	0.2571E-03	0.2056E-03	0.1003E-04	-0.8834E-04	-0.2904E-03	
9	0.1471E-01	0.1516E-02	-0.1989E-03	-0.6349E-04	-0.3765E-04	-0.4146E-04
10	-0.4048E-04	-0.1577E-04	-0.1071E-04	-0.1102E-03	-0.3073E-03	
11	0.1469E-01	0.1514E-02	-0.1477E-03	-0.2748E-04	0.1097E-05	-0.1610E-05
12	-0.8760E-06	0.4645E-05	-0.1337E-04	-0.1398E-03	-0.3259E-03	
13	0.1467E-01	0.1512E-02	-0.1473E-03	-0.2737E-04	0.1046E-05	-0.1768E-05
14	-0.7599E-06	0.4360E-05	-0.1497E-04	-0.1755E-03	-0.3488E-03	
15	0.1463E-01	0.2572E-02	-0.1301E-03	-0.3041E-04	0.1015E-05	-0.1942E-05
16	-0.5834E-06	-0.3737E-05	0.2181E-04	0.7209E-03	-0.3767E-03	
17	0.9521E-02	-0.3709E-02	-0.2313E-03	-0.3028E-05	0.1015E-05	-0.2027E-05
18	-0.4366E-06	0.6026E-04	-0.2027E-03	-0.4348E-02	0.4329E-05	
19	0.4156E-01	0.9738E-02	-0.1277E-03	-0.5392E-04	0.2619E-05	-0.7882E-06
20	0.1422E-05	-0.8433E-04	0.1705E-03	0.4009E-02	0.1934E-02	
21	-0.6397E-02	-0.1535E-02	-0.1815E-03	-0.6359E-04	-0.2989E-04	-0.3448E-04
22	-0.4874E-04	-0.1606E-03	0.1530E-03	0.1183E-01	-0.3044E-02	
23	0.2172E-02	-0.2373E-02	0.1469E-03	0.2149E-03	0.2144E-03	0.2350E-04
24	0.2535E-03	0.4868E-03	-0.8950E-03	-0.1305E-01	-0.7513E-02	
25	0.5965E-02	-0.7227E-02	0.5091E-04	-0.2869E-04	-0.1503E-03	0.3797E-03
26	0.1543E-03	0.4529E-03	-0.1070E-02	0.3084E-02	0.1165E-01	
27	-0.2390E-02	0.1653E-02	-0.5041E-04	0.8774E-03	-0.1592E-02	0.5425E-03
28	-0.6107E-02	0.4016E-02	-0.1141E-02	-0.1061E-01	-0.3689E-01	
29	0.4199E-03	-0.3556E-02	-0.2231E-02	-0.2537E-02	-0.4217E-02	0.1216E-02
30	-0.5153E-02	0.5683E-02	-0.2636E-01	0.1215E-01	-0.1085E-01	

STRESS SIGMAX-1 ON THE UPPER SURFACE (PSI)

	1	2	3	4	5	6
1	0.37353E+06	-0.93206E+06	-0.37420E+07	0.53527E+06	-0.25403E+07	-0.33015E+07
2	-0.35221E+07	0.18032E+07	-0.18433E+07	0.10229E+07	-0.79931E+06	0.78401E+06
3	-0.11745E+07	0.20865E+06	0.16847E+06	0.83808E+06	0.45601E+06	0.78401E+06
4	0.27991E+06	0.36959E+06	-0.16830E+06	0.27924E+06	-0.60414E+05	-0.21918E+06
5	-0.14348E+06	-0.11011E+06	-0.14466E+06	-0.10630E+06	-0.15343E+06	-0.21918E+06
6	-0.13569E+06	-0.25815E+05	0.58318E+05	-0.86936E+04	0.25329E+05	0.51758E+05
7	-0.34529E+06	0.15549E+05	0.62620E+05	0.54744E+05	-0.21952E+05	0.51758E+05
8	0.57341E+05	0.43414E+05	0.31620E+05	0.44659E+05	0.35895E+05	0.43945E+05
9	-0.44958E+06	-0.33459E+06	0.18280E+05	0.23392E+05	0.42455E+05	0.43945E+05
10	0.40451E+05	-0.31975E+06	-0.95450E+05	-0.10392E+05	-0.68640E+05	0.43945E+05
11	0.11360E+07	0.12020E+07	0.43489E+05	0.21110E+05	0.23436E+05	0.23470E+05
12	0.24382E+05	0.10228E+06	0.67373E+05	0.42967E+05	0.52668E+05	0.23470E+05
13	-0.61214E+07	-0.29536E+07	-0.35621E+04	0.36725E+05	0.26542E+05	0.26901E+05
14	0.27841E+05	-0.67614E+04	0.28094E+05	0.34151E+05	0.26952E+05	0.26901E+05
15	0.11277E+08	0.37948E+07	0.64022E+05	0.26001E+05	0.27508E+05	0.27990E+05
16	0.29262E+05	0.39203E+05	0.33896E+05	0.75834E+05	0.14551E+06	0.27990E+05
17	-0.12000E+08	-0.26631E+07	-0.20308E+05	-0.94079E+04	0.28200E+05	0.29734E+05
18	0.30412E+05	0.36157E+05	0.28836E+05	-0.33496E+06	-0.49672E+06	0.29734E+05
19	0.38476E+07	0.10822E+07	0.92648E+05	0.13236E+06	0.28604E+05	0.29877E+05
20	0.32800E+05	0.18610E+05	0.84914E+05	0.14343E+07	0.69651E+06	0.29877E+05
21	0.11474E+07	-0.18167E+07	-0.63829E+05	-0.23160E+06	0.43726E+05	0.36249E+05
22	0.50219E+05	0.67489E+05	-0.64774E+05	-0.17795E+07	0.88031E+06	0.36249E+05
23	0.12780E+08	0.36909E+07	0.73086E+05	0.70521E+05	0.69198E+05	0.10781E+06
24	-0.44415E+06	0.10516E+06	0.95065E+05	-0.15378E+07	-0.90828E+06	0.10781E+06
25	-0.81582E+07	-0.17179E+07	-0.99647E+04	-0.81878E+05	0.36132E+04	-0.74086E+05
26	0.18807E+06	-0.12389E+06	0.28464E+06	0.56125E+07	0.21793E+07	-0.74086E+05
27	0.76467E+06	0.18064E+07	0.40952E+06	-0.39677E+06	0.97282E+06	-0.85492E+06
28	-0.95451E+05	-0.72513E+06	0.12957E+07	-0.70654E+07	-0.65875E+07	-0.85492E+06
29	0.19796E+07	-0.26612E+07	0.29991E+06	-0.12943E+07	0.14330E+07	-0.27625E+07
30	0.56079E+07	-0.14325E+07	0.59477E+06	0.15606E+07	0.17944E+08	-0.27625E+07

STRESS SIGMAY-1 ON THE UPPER SURFACE (PSI)

	1	2	3	4	5	6
1	-0.27213E+06	0.12712E+05	0.29479E+07	-0.28284E+07	0.83921E+06	0.77111E+05
2	0.47547E+07	-0.37667E+07	0.29523E+07	-0.17276E+07	0.10168E+07	-0.71476E+06
3	0.10790E+07	-0.17755E+06	0.23632E+06	-0.60706E+06	0.22395E+06	
4	0.32196E+06	-0.63376E+06	0.49832E+06	-0.40831E+06	0.23171E+06	
5	0.35268E+06	0.13501E+06	0.72985E+05	0.35486E+05	0.68123E+05	0.14133E+06
6	0.94491E+05	0.44631E+05	-0.51986E+05	0.37999E+05	-0.24100E+04	
7	0.52978E+06	0.65478E+05	-0.15264E+05	-0.76776E+04	0.11100E+06	0.39258E+04
8	-0.15018E+05	-0.36038E+04	-0.12736E+05	0.50896E+04	-0.10498E+05	
9	0.54277E+06	0.21438E+06	0.15043E+04	0.97663E+03	-0.55429E+04	-0.59253E+04
10	-0.42804E+04	0.13207E+06	0.44658E+05	0.17196E+05	0.40500E+05	
11	-0.74891E+05	-0.41149E+06	-0.84115E+04	0.19400E+04	0.20295E+04	0.27285E+04
12	0.27565E+04	-0.27454E+05	-0.12636E+05	-0.66187E+04	-0.58205E+04	
13	0.27439E+07	0.13472E+07	0.12656E+05	-0.44576E+04	0.11234E+04	0.17963E+04
14	0.18696E+04	0.15597E+05	0.16377E+04	-0.24461E+04	0.84720E+04	
15	-0.43004E+07	-0.16996E+07	-0.20577E+05	0.15247E+04	0.11862E+04	0.19398E+04
16	0.16135E+04	-0.19177E+04	0.13436E+04	-0.15046E+05	-0.96078E+05	
17	0.62392E+07	0.14450E+07	0.17582E+05	0.14856E+05	0.18690E+04	0.18600E+04
18	0.37516E+04	0.86844E+03	0.24868E+02	0.16336E+06	0.45081E+06	
19	-0.20981E+07	-0.36293E+06	-0.24975E+05	-0.42406E+05	0.13987E+04	0.29586E+04
20	0.13687E+04	0.97676E+04	-0.19454E+05	-0.70129E+06	-0.27476E+06	
21	-0.35346E+06	0.70788E+06	0.36543E+05	0.99957E+05	-0.25547E+04	0.10625E+04
22	-0.68034E+04	-0.13307E+05	0.67754E+05	0.98444E+06	-0.13468E+07	
23	-0.41729E+07	-0.14816E+07	-0.53363E+04	-0.36603E+05	-0.77468E+04	-0.95955E+05
24	0.69530E+06	-0.46651E+05	-0.18924E+05	0.10311E+07	0.45486E+06	
25	0.44767E+07	0.31618E+06	-0.12787E+05	0.11500E+06	-0.78627E+05	0.21140E+06
26	-0.26649E+06	0.12142E+06	-0.19560E+06	-0.26191E+07	-0.57021E+06	
27	-0.13013E+07	-0.13364E+06	-0.58818E+06	0.92387E+06	-0.15574E+07	0.17338E+07
28	-0.84588E+06	0.95559E+06	-0.15501E+07	0.39347E+07	0.28623E+07	
29	-0.10184E+06	0.96416E+06	-0.48291E+06	0.13914E+07	-0.21393E+07	0.39894E+07
30	-0.71536E+07	0.18221E+07	-0.18171E+07	0.19797E+07	-0.12299E+08	

## STRESS SIGMAXY-1 ON THE UPPER SURFACE (PSI)

	1	2	3	4	5	6
1	0.67555E+05	0.50684E+06	-0.30073E+06	0.25275E+07	0.93369E+06	0.20245E+07
2	-0.18543E+07	0.24413E+07	-0.15319E+07	0.96850E+06	-0.44567E+06	0.13346E+06
3	-0.18696E+06	0.11152E+06	-0.37184E+06	0.81258E+05	-0.48728E+06	0.20429E+05
4	-0.47931E+06	0.32282E+06	-0.33786E+06	0.20619E+06	-0.16655E+06	-0.28801E+05
5	-0.25158E+06	-0.72370E+05	0.44134E+05	0.48041E+05	0.45892E+05	-0.13493E+05
6	0.17764E+05	-0.46894E+04	0.22826E+05	-0.22198E+05	-0.32808E+04	-0.74285E+04
7	-0.27930E+06	-0.59446E+05	-0.19291E+05	-0.19116E+05	-0.83413E+05	-0.84764E+04
8	-0.14950E+05	-0.16090E+05	-0.19012E+03	-0.23454E+05	-0.22417E+04	-0.88370E+04
9	-0.21719E+06	0.29897E+05	-0.31119E+04	-0.62986E+04	-0.12827E+05	-0.92883E+04
10	-0.12604E+05	0.99456E+05	0.32101E+05	0.41657E+04	0.20486E+05	-0.96044E+04
11	-0.70312E+06	-0.42386E+06	-0.11050E+05	-0.54405E+04	-0.69655E+04	-0.11218E+05
12	-0.77555E+04	-0.31219E+05	-0.20620E+05	-0.10132E+05	-0.17312E+05	0.35165E+05
13	0.15165E+07	0.77468E+06	0.20130E+04	-0.99248E+04	-0.79078E+04	-0.13422E+06
14	-0.87799E+04	0.21077E+04	-0.77353E+04	-0.67731E+04	-0.12494E+05	-0.98150E+06
15	-0.36037E+07	-0.10943E+07	-0.15547E+05	-0.72869E+04	-0.81804E+04	-0.20800E+07
16	-0.89423E+04	-0.11534E+05	-0.11137E+05	-0.34614E+05	0.27803E+05	0.11280E+07
17	0.25246E+07	0.66982E+06	0.10990E+05	0.30888E+04	-0.88857E+04	-0.18837E+06
18	-0.10820E+05	-0.12216E+05	0.18914E+03	0.15872E+06	-0.19005E+06	0.21959E+07
19	-0.10936E+07	-0.53524E+06	-0.28765E+05	-0.37100E+05	-0.80316E+04	0.84562E+06
20	-0.88928E+04	-0.42803E+04	-0.32111E+05	-0.41627E+06	-0.31831E+05	0.21959E+07
21	-0.34919E+06	0.58132E+06	0.20604E+05	0.73386E+05	-0.13418E+05	0.11280E+07
22	-0.12791E+05	-0.16664E+05	-0.16358E+04	0.20845E+06	0.92714E+06	-0.20800E+07
23	-0.44700E+07	-0.10543E+07	-0.32129E+05	0.78003E+03	-0.21667E+05	0.98150E+06
24	-0.37948E+06	-0.11295E+05	-0.67496E+03	0.43229E+06	-0.26496E+06	-0.13422E+06
25	0.13428E+07	0.83641E+06	0.34573E+05	-0.37565E+05	0.78977E+05	0.35165E+05
26	0.15725E+06	-0.14916E+05	0.13795E+05	-0.16229E+07	-0.76023E+06	-0.13422E+06
27	0.44438E+06	-0.96481E+06	0.29000E+06	-0.59136E+06	0.84562E+06	-0.98150E+06
28	0.77984E+06	-0.59119E+06	0.76595E+06	0.12914E+07	0.21959E+07	-0.20800E+07
29	-0.10513E+07	0.81948E+06	0.31737E+06	-0.47745E+06	0.11280E+07	0.20800E+07
30	0.31680E+07	-0.74510E+06	0.17796E+07	-0.38980E+07	-0.18837E+06	0.20800E+07

STRESS SIGMAX ON THE UPPER SURFACE (PSI)

	1	2	3	4	5	6
1	0.37353E+06	-0.93206E+06	-0.37420E+07	0.53527E+06	-0.25403E+07	-0.33015E+07
2	-0.35221E+07	0.18032E+07	-0.18433E+07	0.10229E+07	-0.79931E+06	0.78401E+06
3	-0.11745E+07	0.20865E+06	0.16847E+06	0.83808E+06	0.45601E+06	0.78401E+06
4	0.27991E+06	0.36959E+06	-0.16830E+06	0.27924E+06	-0.60414E+05	-0.21931E+06
5	-0.14348E+06	-0.11016E+06	-0.14457E+06	-0.10643E+06	-0.15334E+06	-0.21931E+06
6	-0.13558E+06	-0.25868E+05	0.58388E+05	-0.86698E+04	0.25263E+05	0.51938E+05
7	-0.34536E+06	0.15665E+05	0.62453E+05	0.54924E+05	-0.22119E+05	0.51938E+05
8	0.57146E+05	0.43633E+05	0.31609E+05	0.44686E+05	0.35800E+05	0.44027E+05
9	-0.44961E+06	-0.33470E+06	0.18494E+05	0.23126E+05	0.42384E+05	0.44027E+05
10	0.40394E+05	-0.31996E+06	-0.95232E+05	-0.10425E+05	-0.68742E+05	0.23374E+05
11	0.11359E+07	0.12021E+07	0.43291E+05	0.21406E+05	0.23514E+05	0.23374E+05
12	0.24485E+05	0.10268E+06	0.67207E+05	0.43085E+05	0.52535E+05	0.27001E+05
13	-0.61214E+07	-0.29537E+07	-0.33026E+04	0.36411E+05	0.26458E+05	0.27001E+05
14	0.27739E+05	-0.71401E+04	0.28509E+05	0.34039E+05	0.26854E+05	0.27884E+05
15	0.11277E+08	0.37949E+07	0.63839E+05	0.26301E+05	0.27587E+05	0.27884E+05
16	0.29386E+05	0.39685E+05	0.33607E+05	0.76081E+05	0.14533E+06	0.29822E+05
17	-0.12000E+08	-0.26632E+07	-0.20110E+05	-0.96578E+04	0.28133E+05	0.29822E+05
18	0.30304E+05	0.35752E+05	0.29349E+05	-0.33514E+06	-0.49681E+06	0.29799E+05
19	0.38475E+07	0.10822E+07	0.92530E+05	0.13259E+06	0.28664E+05	0.29799E+05
20	0.32896E+05	0.18983E+05	0.84625E+05	0.14347E+07	0.69627E+06	0.36307E+05
21	0.11474E+07	-0.18168E+07	-0.63734E+05	-0.23174E+06	0.43690E+05	0.36307E+05
22	0.50145E+05	0.67216E+05	-0.64377E+05	-0.17796E+07	0.88022E+06	0.10799E+06
23	0.12780E+08	0.36909E+07	0.73048E+05	0.70636E+05	0.69081E+05	0.10799E+06
24	-0.44434E+06	0.10534E+06	0.94872E+05	-0.15375E+07	-0.90847E+06	-0.74178E+05
25	-0.81583E+07	-0.17179E+07	-0.99349E+04	-0.81910E+05	0.36915E+04	-0.74178E+05
26	0.18820E+06	-0.12399E+06	0.28476E+06	0.56125E+07	0.21793E+07	-0.85492E+06
27	0.76467E+06	0.18064E+07	0.40952E+06	-0.39677E+06	0.97282E+06	-0.85492E+06
28	-0.95451E+05	-0.72513E+06	0.12957E+07	-0.70654E+07	-0.65875E+07	-0.27625E+07
29	0.19796E+07	-0.26612E+07	0.29991E+06	-0.12943E+07	0.14330E+07	-0.27625E+07
30	0.56079E+07	-0.14325E+07	0.59477E+06	0.15606E+07	0.17944E+08	-0.27625E+07

STRESS SIGMAY ON THE UPPER SURFACE (PSI)

	1	2	3	4	5	6
1	-0.27213E+06	0.12712E+05	0.29479E+07	-0.28284E+07	0.83920E+06	0.77007E+05
2	0.47547E+07	-0.37688E+07	0.29523E+07	-0.17276E+07	0.10168E+07	-0.71465E+06
3	0.10790E+07	-0.17755E+06	0.23624E+06	-0.60699E+06	0.22382E+06	0.14115E+06
4	0.32187E+06	-0.63370E+06	0.49836E+06	-0.40831E+06	0.23171E+06	0.42849E+04
5	0.35268E+06	0.13501E+06	0.73104E+05	0.35332E+05	0.68332E+05	-0.58325E+04
6	0.94739E+05	0.44611E+05	-0.51964E+05	0.37999E+05	-0.24100E+04	0.26328E+04
7	0.52978E+06	0.65478E+05	-0.15389E+05	-0.74491E+04	0.11079E+06	0.18958E+04
8	-0.15276E+05	-0.34390E+04	-0.12759E+05	0.50896E+04	-0.10498E+05	0.18416E+04
9	0.54277E+06	0.21438E+06	0.16565E+04	0.77033E+03	-0.56005E+04	0.19463E+04
10	-0.43514E+04	0.13205E+06	0.44680E+05	0.17196E+05	0.40500E+05	0.28863E+04
11	-0.74891E+05	-0.41149E+06	-0.85276E+04	0.21797E+04	0.21090E+04	0.11219E+04
12	0.28604E+04	-0.27289E+05	-0.12659E+05	-0.66187E+04	-0.58205E+04	-0.95726E+05
13	0.27439E+07	0.13472E+07	0.12749E+05	-0.46815E+04	0.10424E+04	0.21131E+06
14	0.17482E+04	0.15345E+05	0.17023E+04	-0.24461E+04	0.84720E+04	0.17339E+07
15	-0.43004E+07	-0.16996E+07	-0.20650E+05	0.17078E+04	0.12621E+04	0.39893E+07
	0.17394E+04	-0.15636E+04	0.12262E+04	-0.15046E+05	-0.96078E+05	0.28623E+07
	0.62392E+07	0.14450E+07	0.17626E+05	0.14720E+05	0.17999E+04	-0.21393E+07
	0.36347E+04	0.44239E+03	0.20239E+03	0.16336E+06	0.45081E+06	0.12299E+08
	-0.20981E+07	-0.36293E+06	-0.24997E+05	-0.42317E+05	0.14561E+04	
	0.14666E+04	0.10229E+05	-0.19695E+05	-0.70129E+06	-0.27476E+06	
	-0.35346E+06	0.70788E+06	0.36556E+05	0.99934E+05	-0.25950E+04	
	-0.68684E+04	-0.13719E+05	0.68049E+05	0.98444E+06	-0.13468E+07	
	-0.41729E+07	-0.14816E+07	-0.53583E+04	-0.36514E+05	-0.78938E+04	
	0.69506E+06	-0.46255E+05	-0.19171E+05	0.10311E+07	0.45486E+06	
	0.44767E+07	0.31618E+06	-0.12774E+05	0.11498E+06	-0.78485E+05	
	-0.26628E+06	0.12120E+06	-0.19534E+06	-0.26191E+07	-0.57021E+06	
	-0.13013E+07	-0.13364E+06	-0.58816E+06	0.92391E+06	-0.15574E+07	
	-0.84596E+06	0.95570E+06	-0.15502E+07	0.39347E+07	0.28623E+07	
	-0.10184E+06	0.96416E+06	-0.48294E+06	0.13913E+07	-0.21393E+07	
	-0.71536E+07	0.18220E+07	-0.18171E+07	0.19797E+07	-0.12299E+08	

STRESS SIGMAXY ON THE UPPER SURFACE (PSI)

	1	2	3	4	5	6
1	0.67555E+05	0.50684E+06	-0.30073E+06	0.25275E+07	0.93369E+06	0.20245E+07
2	-0.18543E+07	0.24413E+07	-0.15319E+07	0.96850E+06	-0.44567E+06	0.13337E+06
3	-0.18696E+06	0.11152E+06	-0.37189E+06	0.81297E+05	-0.48724E+06	0.20575E+05
4	-0.47928E+06	0.32279E+06	-0.33782E+06	0.20626E+06	-0.16655E+06	-0.29046E+05
5	-0.25158E+06	-0.72444E+05	0.44227E+05	0.47878E+05	0.45729E+05	-0.13494E+05
6	0.17600E+05	-0.45825E+04	0.22852E+05	-0.22169E+05	-0.32808E+04	-0.74268E+04
7	-0.27930E+06	-0.59401E+05	-0.19470E+05	-0.18916E+05	-0.83212E+05	-0.84778E+04
8	-0.14685E+05	-0.16252E+05	-0.12259E+03	-0.23447E+05	-0.22417E+04	-0.88365E+04
9	-0.21719E+06	0.29817E+05	-0.29382E+04	-0.65364E+04	-0.12825E+05	-0.92873E+04
10	-0.12603E+05	0.99295E+05	0.32169E+05	0.41729E+04	0.20486E+05	-0.96047E+04
11	-0.70312E+06	-0.42379E+06	-0.11235E+05	-0.51680E+04	-0.69645E+04	-0.11215E+05
12	-0.77592E+04	-0.30957E+05	-0.20741E+05	-0.10102E+05	-0.17312E+05	0.34990E+05
13	0.15165E+07	0.77464E+06	0.21840E+04	-0.10197E+05	-0.79075E+04	-0.13409E+06
14	-0.87793E+04	0.17235E+04	-0.75393E+04	-0.68073E+04	-0.12494E+05	-0.98155E+06
15	-0.36037E+07	-0.10942E+07	-0.15692E+05	-0.70505E+04	-0.81808E+04	-0.20800E+07
	-0.89455E+04	-0.11106E+05	-0.11412E+05	-0.34531E+05	0.27803E+05	
	0.25246E+07	0.66979E+06	0.11101E+05	0.28821E+04	-0.88864E+04	
	-0.10819E+05	-0.12690E+05	0.51715E+03	0.15862E+06	-0.19005E+06	
	-0.10936E+07	-0.53522E+06	-0.28833E+05	-0.36954E+05	-0.80320E+04	
	-0.88921E+04	-0.38523E+04	-0.32485E+05	-0.41612E+06	-0.31831E+05	
	-0.34919E+06	0.58131E+06	0.20646E+05	0.73296E+05	-0.13418E+05	
	-0.12789E+05	-0.17020E+05	-0.13038E+04	0.20830E+06	0.92714E+06	
	-0.44700E+07	-0.10543E+07	-0.32086E+05	0.68957E+03	-0.21508E+05	
	-0.37920E+06	-0.11016E+05	-0.98653E+03	0.43239E+06	-0.26496E+06	
	0.13428E+07	0.83642E+06	0.34580E+05	-0.37490E+05	0.78905E+05	
	0.15710E+06	-0.15065E+05	0.13965E+05	-0.16230E+07	-0.76023E+06	
	0.44438E+06	-0.96478E+06	0.29003E+06	-0.59136E+06	0.84566E+06	
	0.77988E+06	-0.59114E+06	0.76587E+06	0.12914E+07	0.21959E+07	
	-0.10513E+07	0.81948E+06	0.31737E+06	-0.47745E+06	0.11280E+07	
	0.31680E+07	-0.74510E+06	0.17796E+07	-0.38980E+07	-0.18837E+06	

## VITA

Prabh Nandan Singh was born on August 26, 1946, in Amritsar, Panjab, India. He attended G.R.D. Khalsa Higher Secondary School at Amritsar and passed in June, 1962. He then attended Panjab Engineering College, Chandigarh, India, and graduated in August, 1967, with a Bachelor of Science degree in Mechanical Engineering. After a few years in industry he was admitted to the Graduate School of Louisiana State University and is now a candidate for the Master of Science degree in Nuclear Engineering in August, 1972.